

Signals and degeneracies of the primordial power spectrum: **Neutrinos & Inflation**

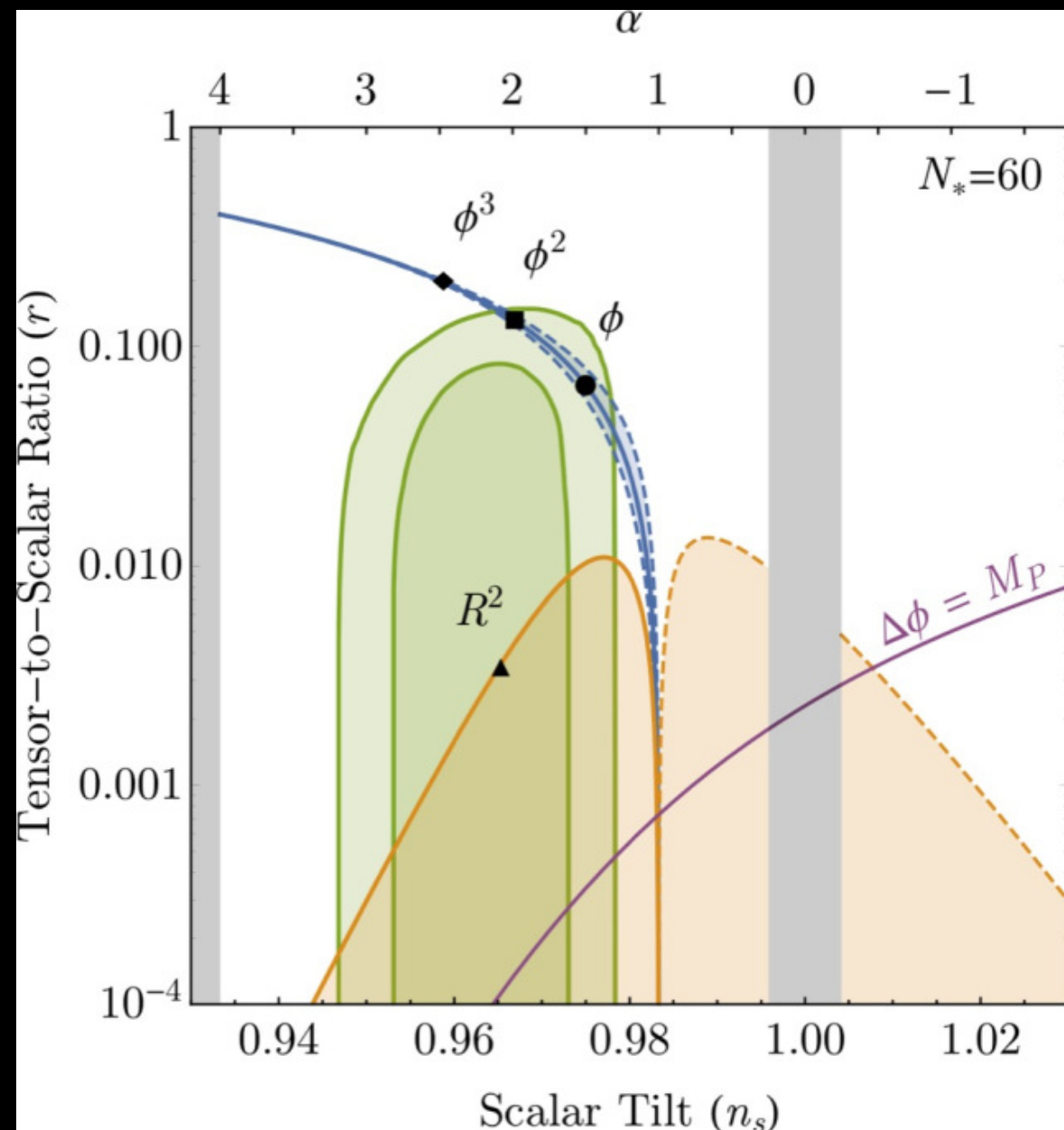
Kevork Abazajian
University of California, Irvine

Cosmology with the CMB and its Polarization

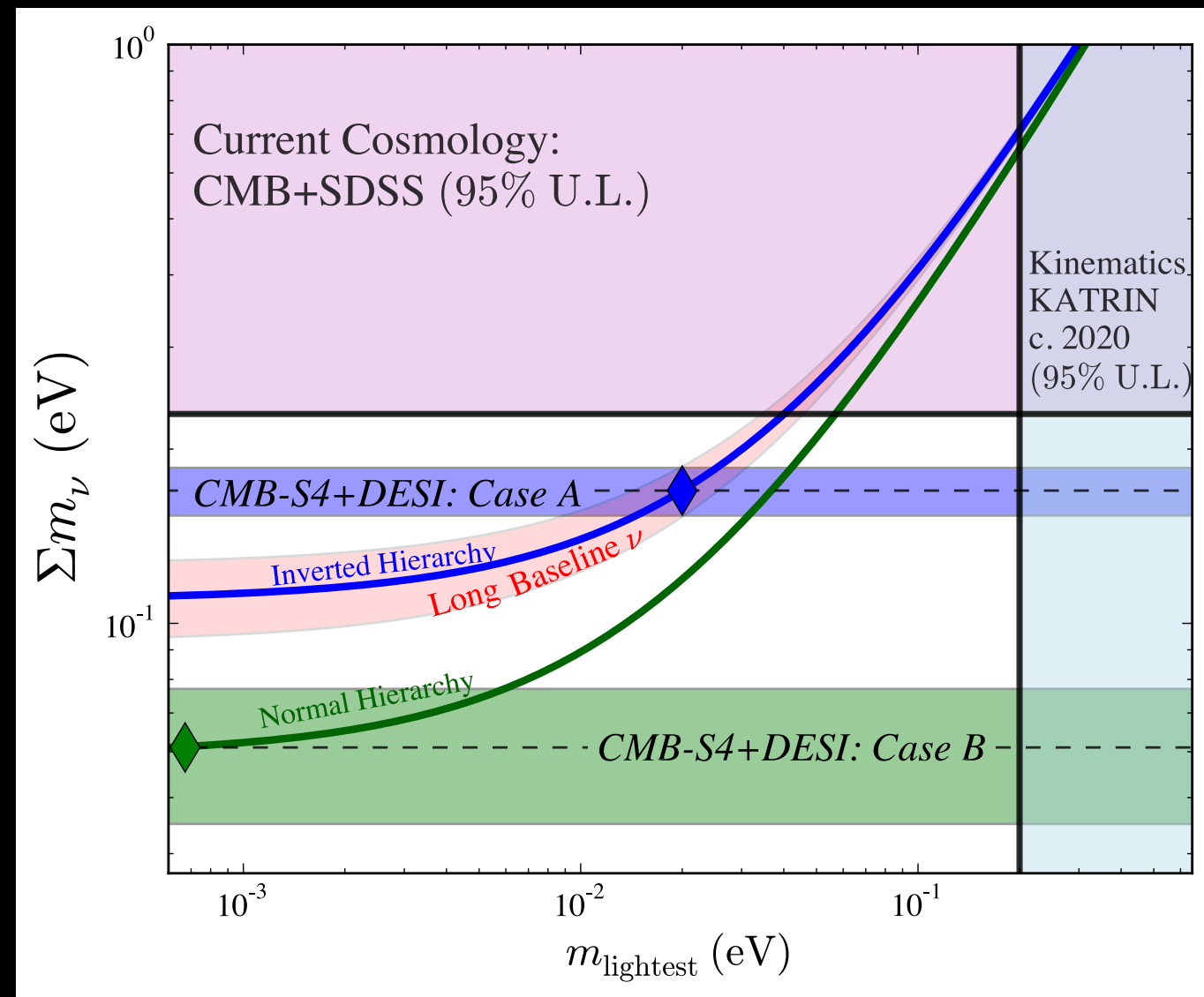
Fine Theoretical Physics Institute
University of Minnesota
January 14, 2015

Two Fundamental Physics Goals of Cosmology

inflationary tilt &
scalar-to-tensor ratio



the absolute
neutrino mass scale



Super-K

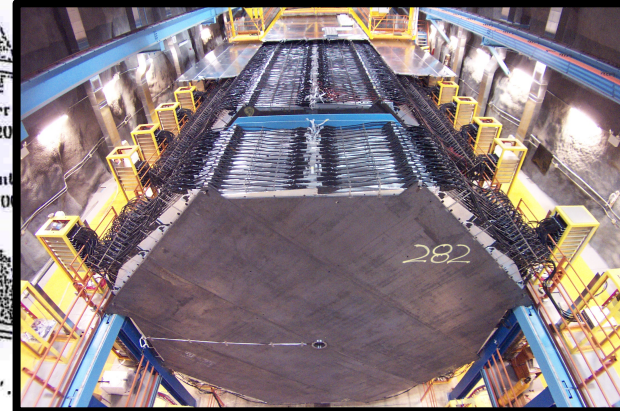
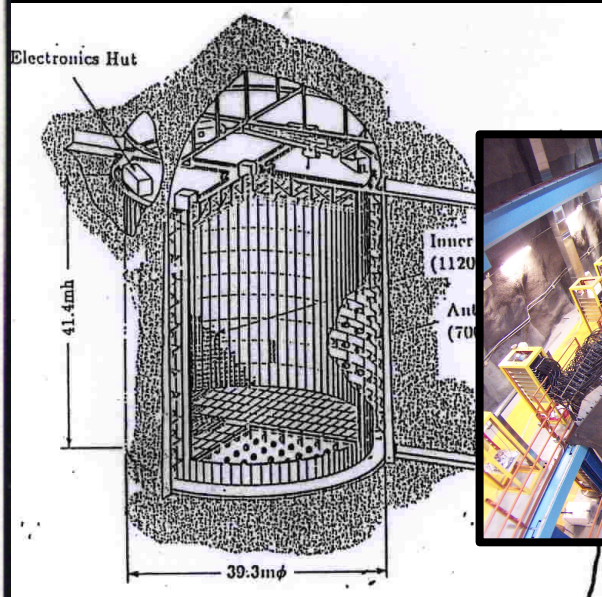
Zenith - Angle
Dependence
of $\nu_\mu/\bar{\nu}_\mu$

$L \sim 10^6 \text{m}$

Earth

Final word:
MINOS!

$$\delta m^2 \sim 3 \times 10^{-3} \text{eV}^2 \text{ Flux}$$



$L \sim 10^{11} \text{m}$

Earth

Solar ν 's

A suppression
of
the expected
 ν_e
Flux

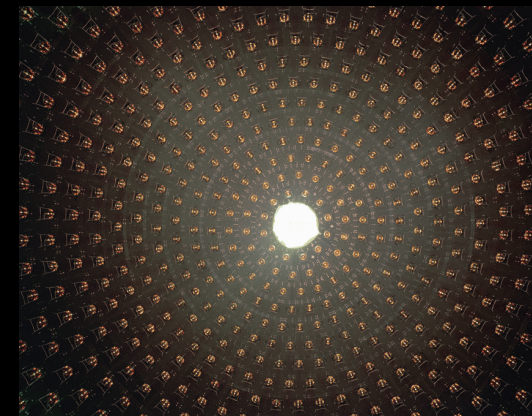
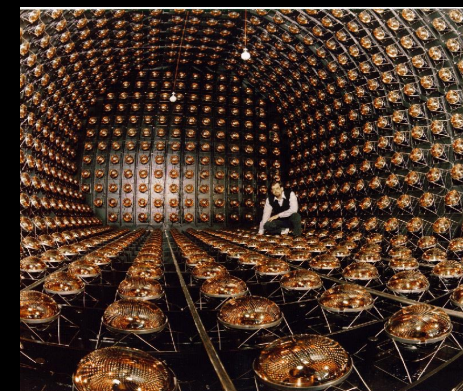
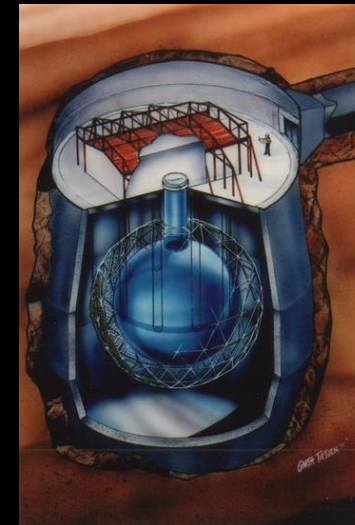
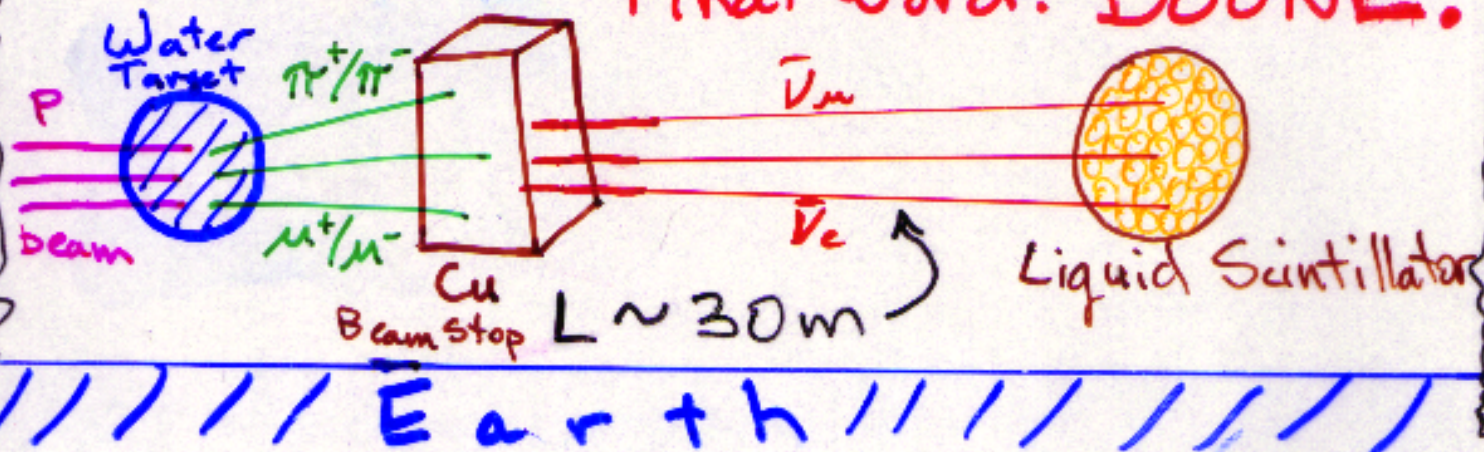
Final word:
KamLAND!!

$$\delta m^2 \sim 10^{-5} \text{eV}^2$$

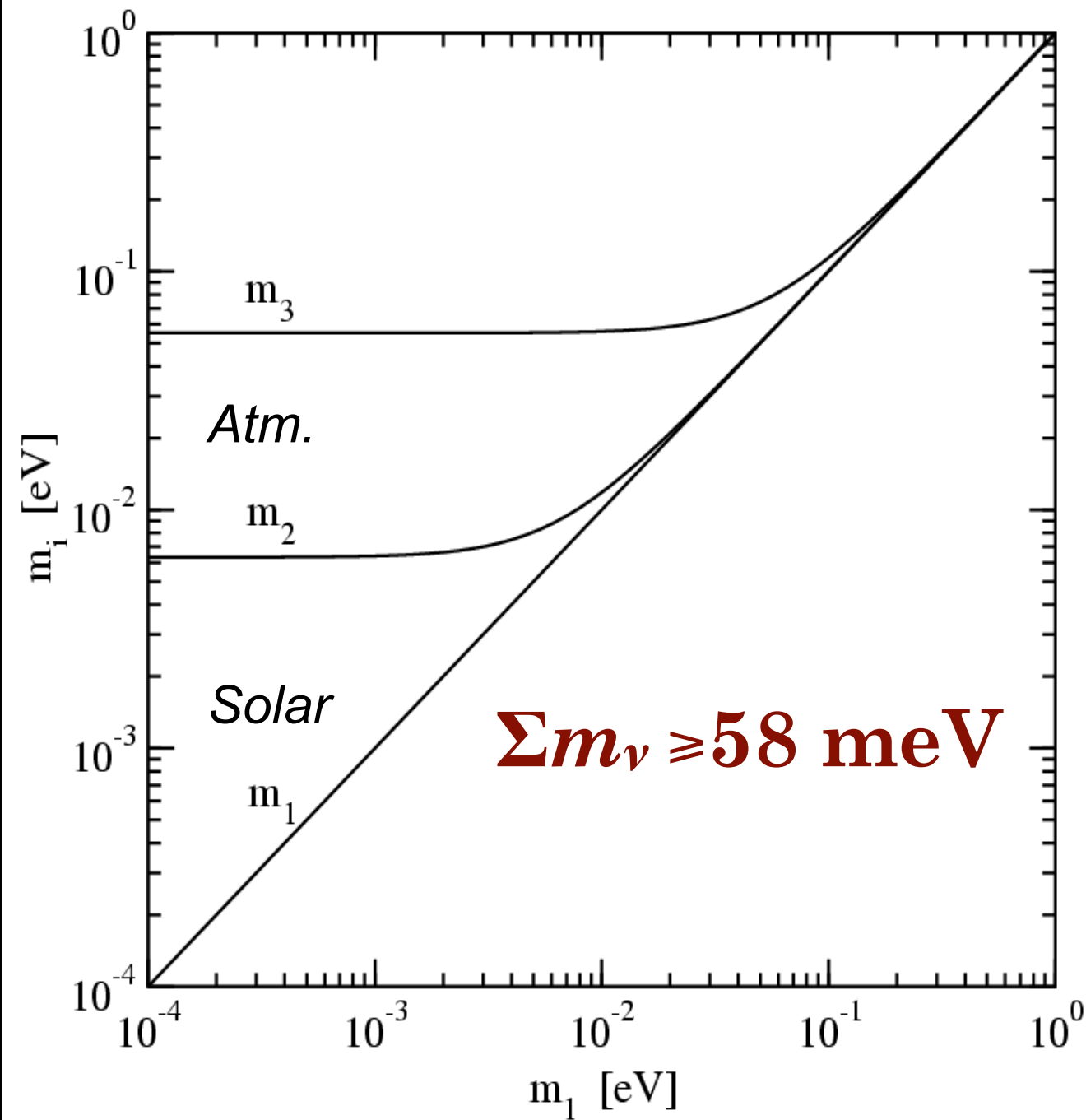
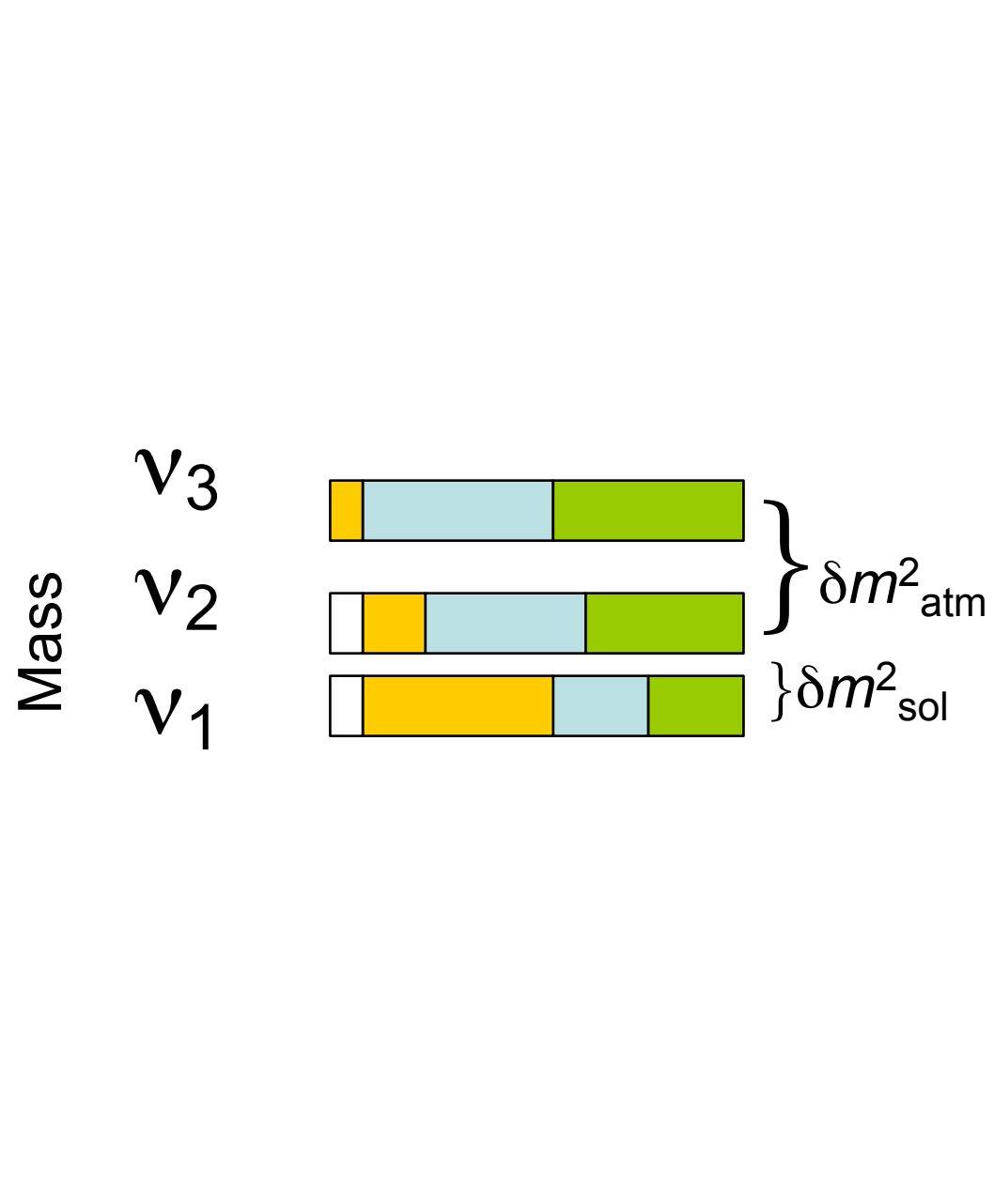
SNO: Appearance!

LSND

$0.1 \text{eV}^2 < \delta m^2 < 6 \text{eV}^2$
Final Word: BOONE!



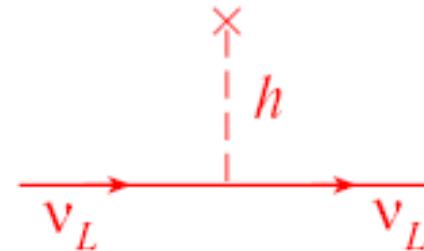
The absolute neutrino mass scale



What is the absolute neutrino mass scale?

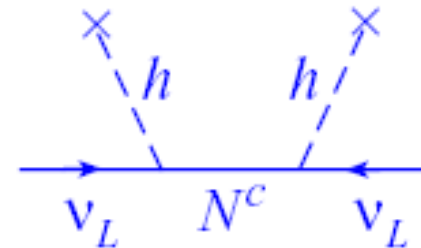
Tree-level mass generation

$(\Delta L = 0)$

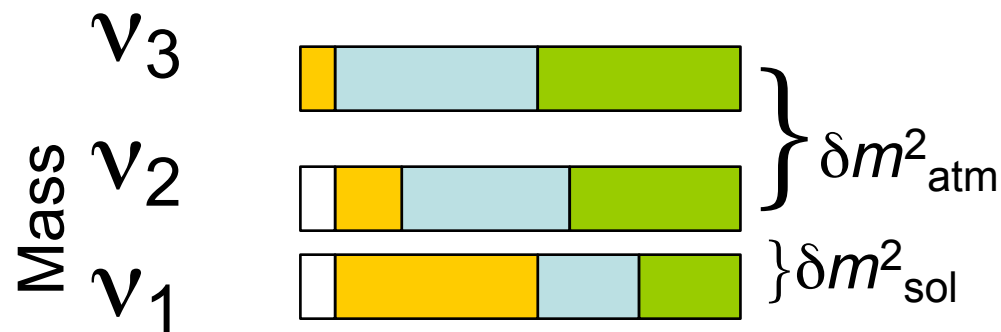


Dirac mass m_D

$(\Delta L = 2)$

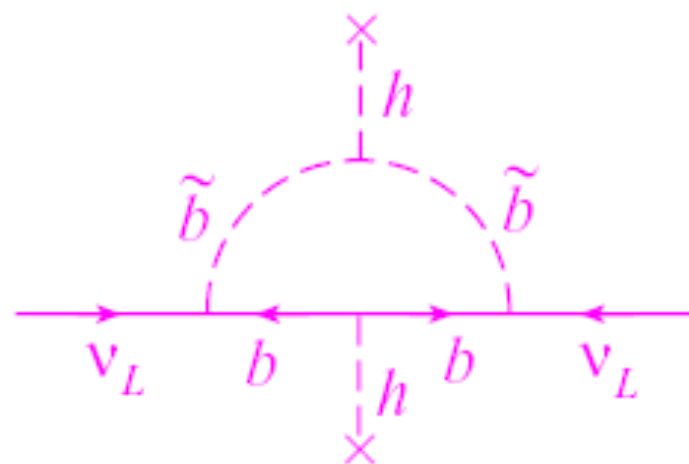


Majorana mass m_M



Radiative mass generation via new interactions

e.g.

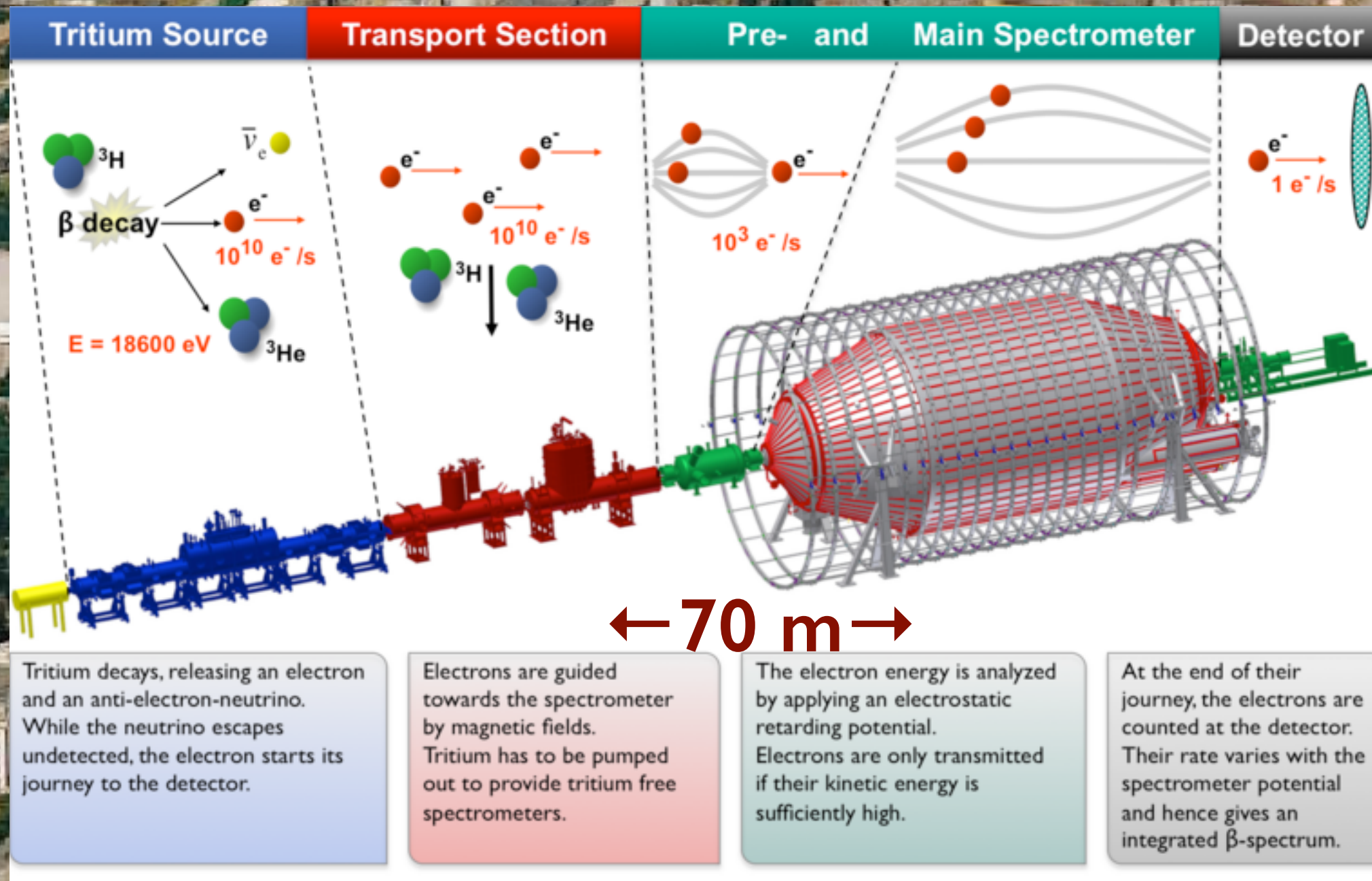


Ma and Sarkar (98)
Drees et al. (98)

Cheung & Kong (99)
Babu (98)
Mohapatra & Senjanovic (81)
Zee (80)

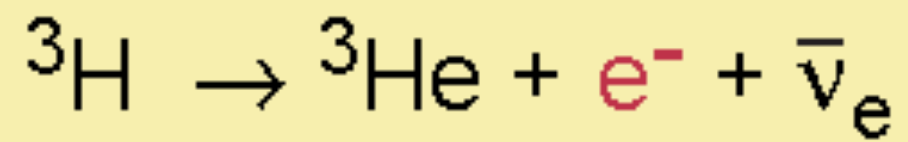
The absolute neutrino mass scale in the laboratory

Forschungszentrum Karlsruhe



KATRIN & TLK

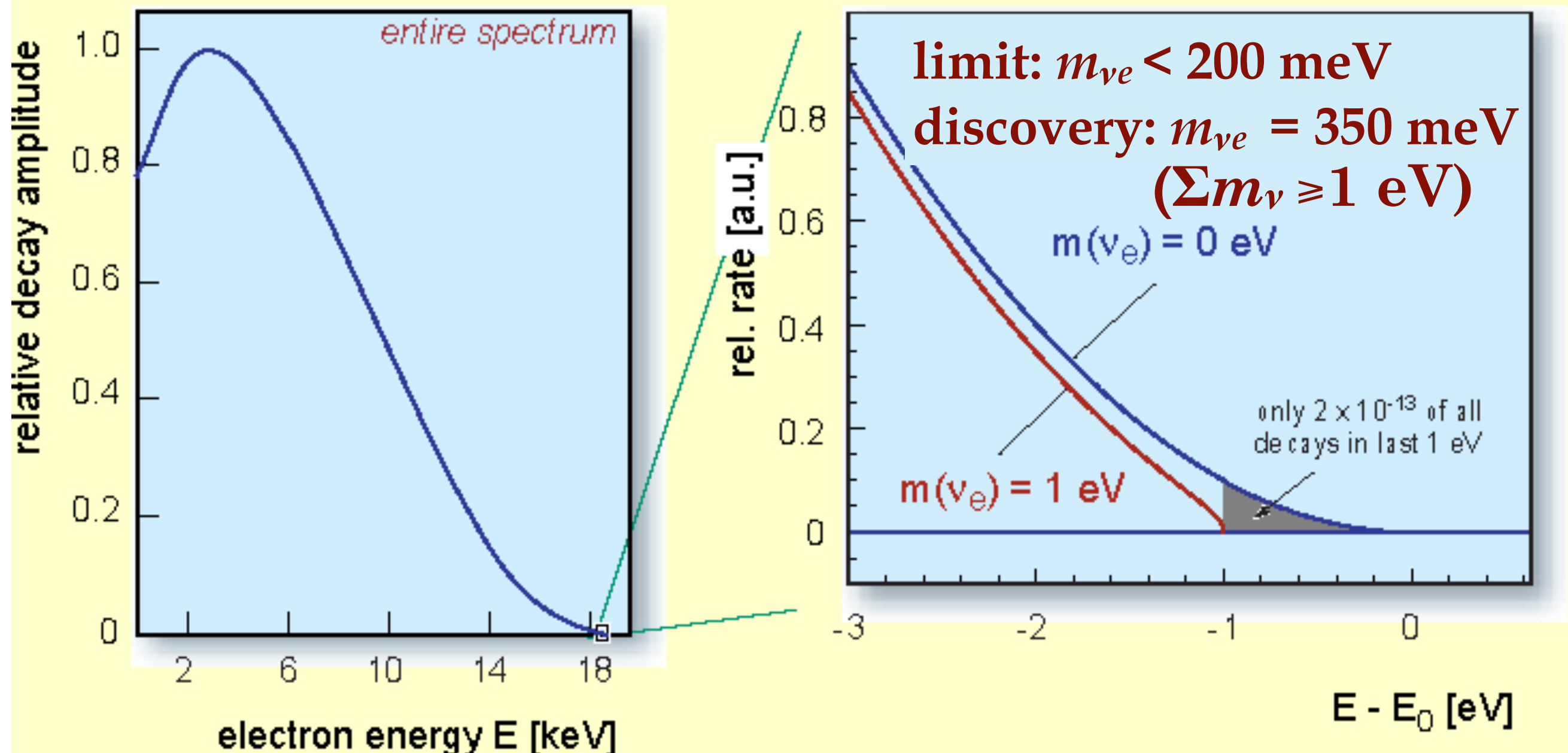
tritium β -decay and the neutrino rest mass



superallowed

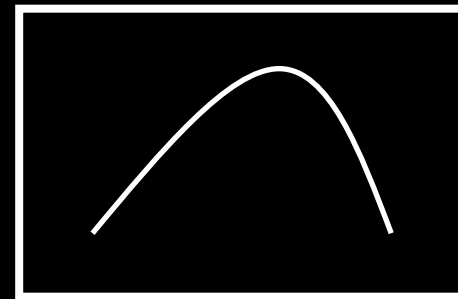
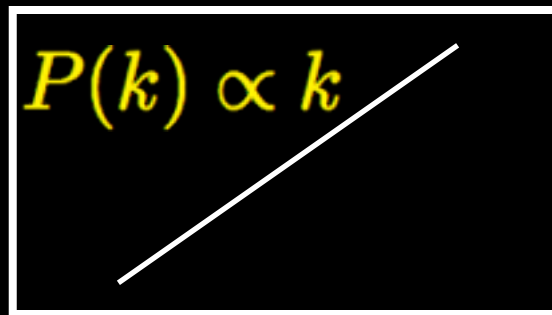
half life : $t_{1/2} = 12.32 \text{ a}$

β end point energy : $E_0 = 18.57 \text{ keV}$



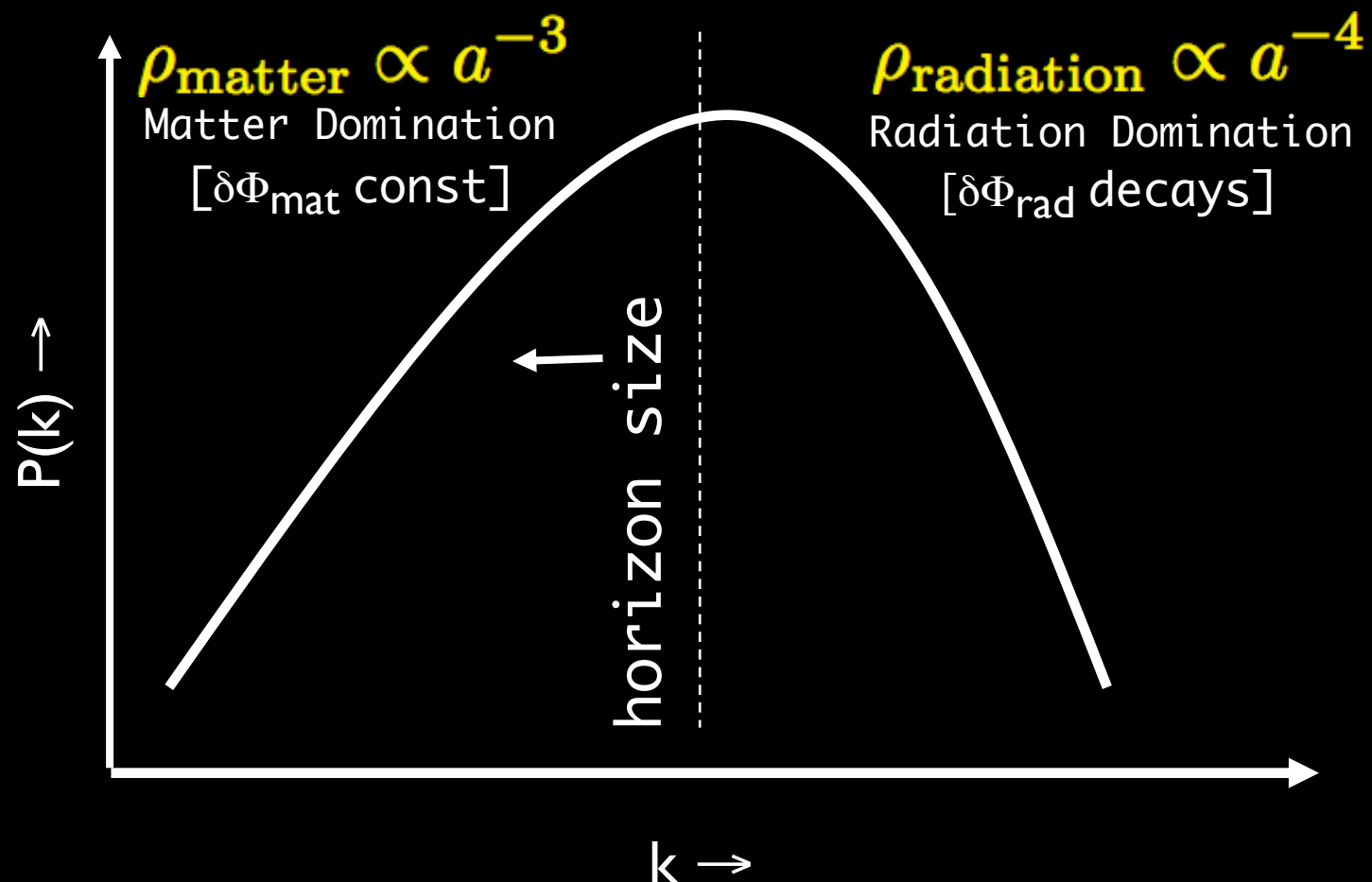
The Cosmological Matter Power Spectrum

Inflation:



?

Perturbations enter horizon:



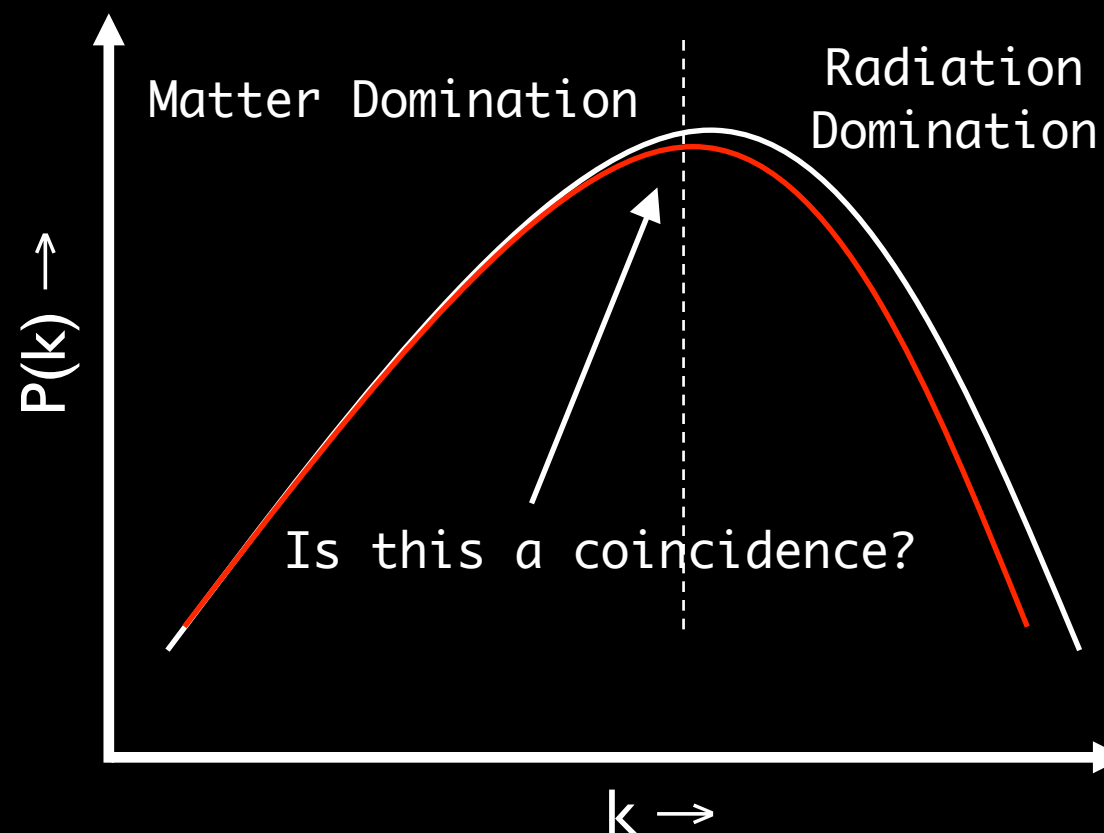
How does probe neutrinos?

$$n_\nu = N_\nu \times \left(\frac{3}{11}\right) n_\gamma \approx 340 \text{ cm}^{-3} \quad (\text{Assuming thermal equilibrium})$$

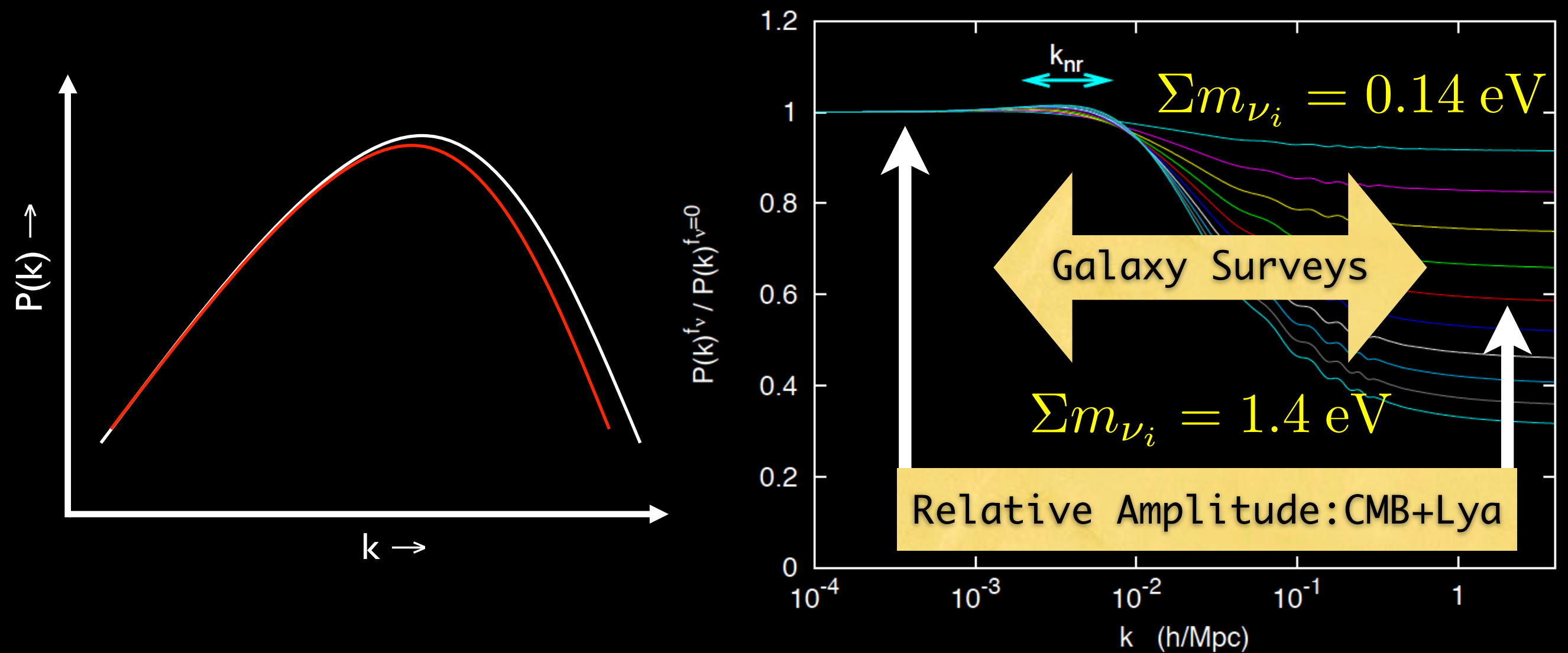
$$\rho_\nu = \sum m_i n_{\nu_i}$$

$$\Omega_\nu \approx \frac{\sum m_{\nu_i}}{93 \text{ h}^2 \text{ eV}}$$

$$E^2 = p^2 + m^2$$

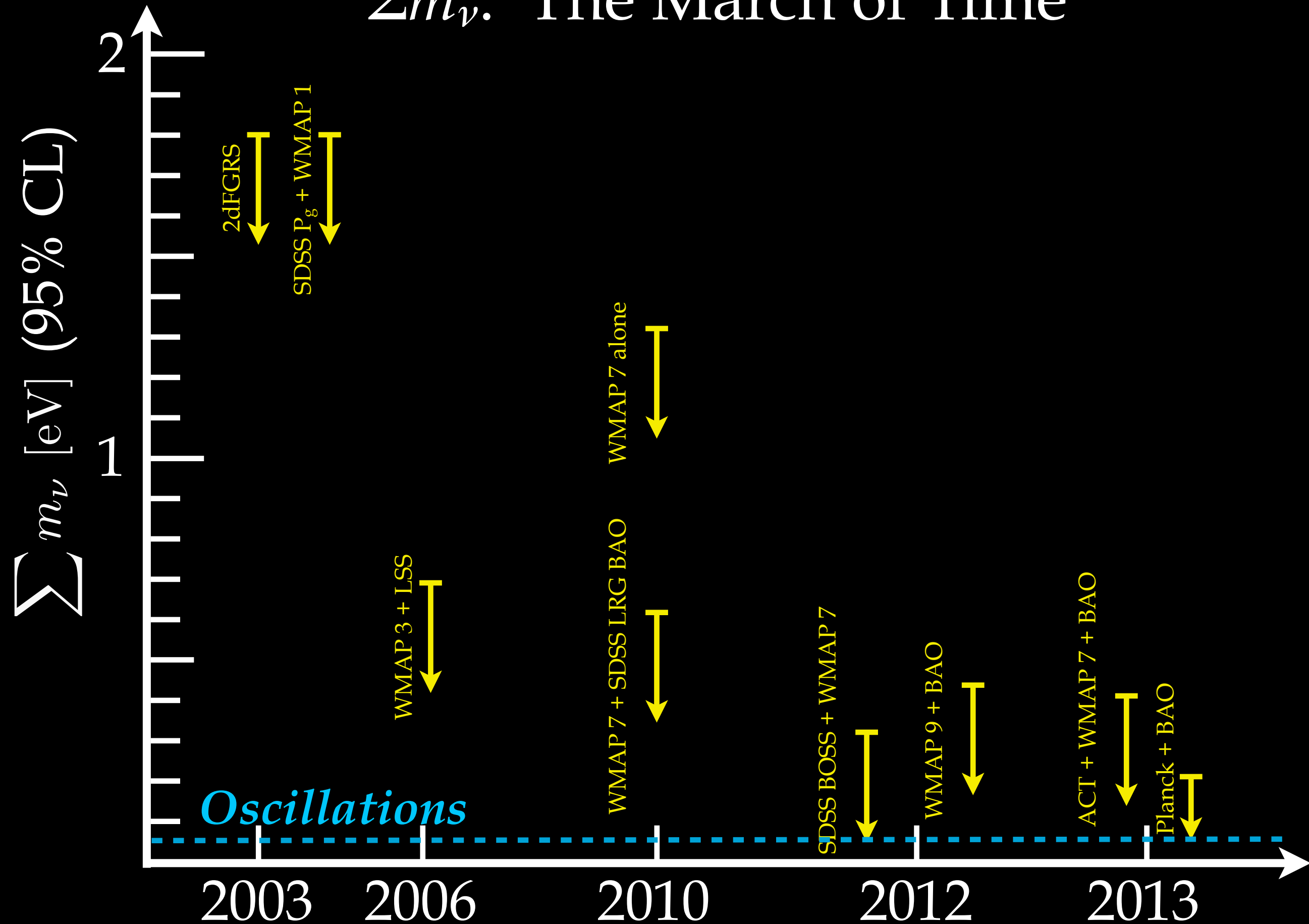


Distinguishing Features in the Power Spectrum

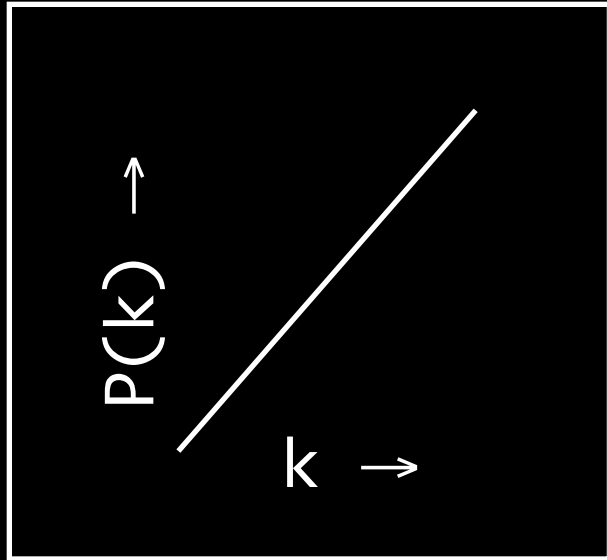


1. Shape Information:
Galaxy Surveys (Future: CMB lensing, Weak Lensing)
2. Relative Amplitude Information:
CMB plus Lyman-alpha Forest, Galaxy Bias $\frac{\Delta P(k)}{P(k)} = -8 \frac{\Omega_\nu}{\Omega_m}$

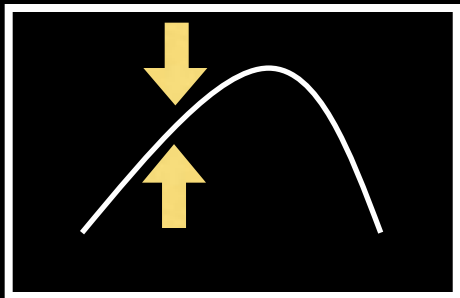
Σm_ν : The March of Time



The Primordial Spectrum: CMB gives a Precision Determination at Large Scales



$$P(k) = Ak^n$$

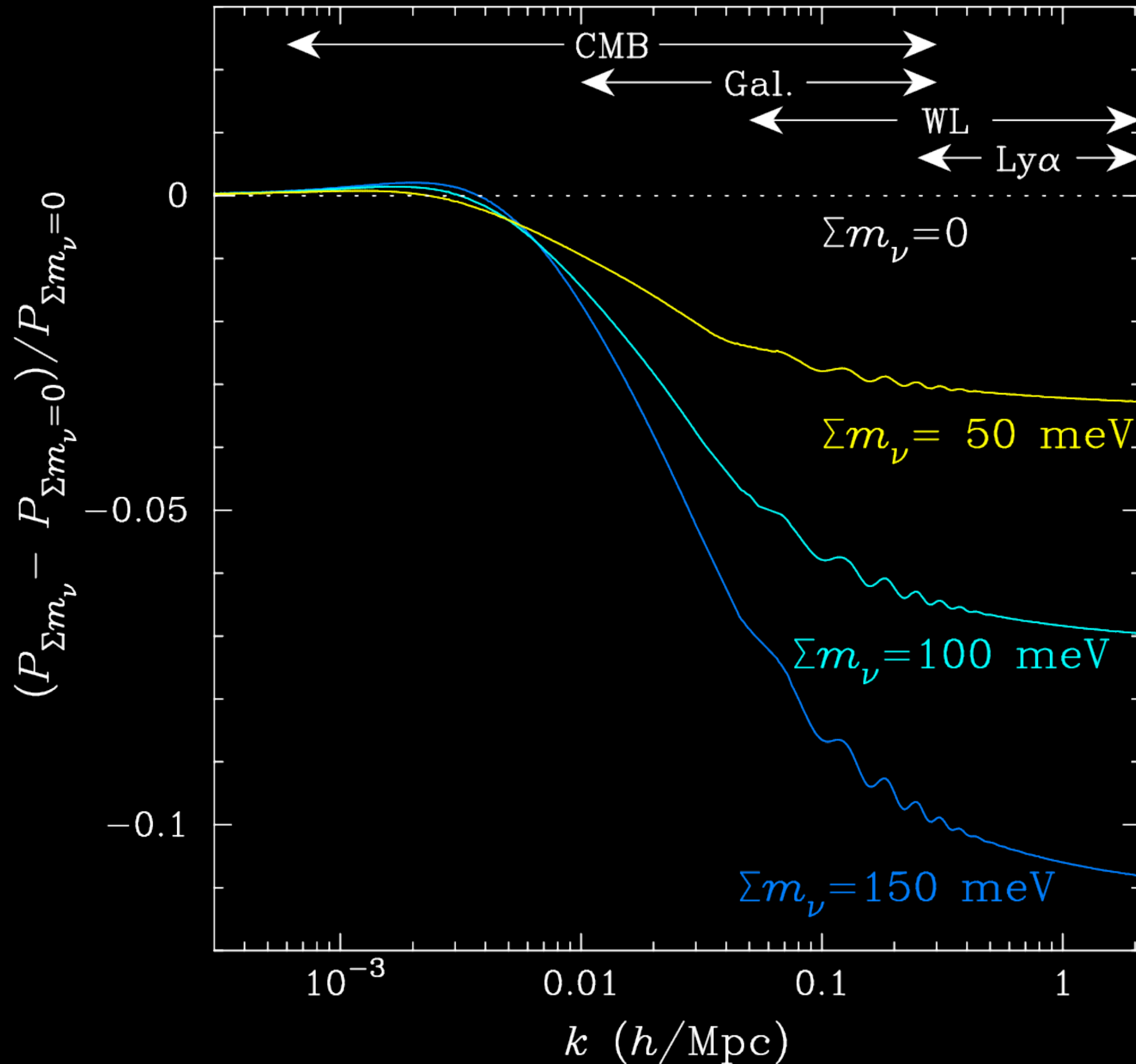


Planck 2013 + WMAP Pol.: (Planck Collab. 2013)

$$A = 2.196^{+0.051}_{-0.060} \quad (3\%)$$

$$n = 0.9603 \pm 0.0073 \quad (0.8\%)$$

Upcoming high- k High-Precision Era: Relative Change to $P(k)$



Estimating Upcoming Cosmological Neutrino Mass Sensitivities

$$\frac{\Delta P(k)}{P(k)} \approx 1\% \approx -8 \frac{\Omega_\nu}{\Omega_m}$$

Hu, Eisenstein & Tegmark 1998

$$\Omega_\nu \approx \frac{\sum m_{\nu_i}}{93 h^2 \text{ eV}}$$

$$\Rightarrow m_\nu \lesssim (1\%/8) \times \Omega_m (93 h^2 \text{ eV})$$

$$\Rightarrow m_\nu \lesssim 20 \text{ meV}$$

Kaplinghat et al PRL 2003 (CMB WL)

Wang et al PRL 2005 (WL Clusters)

De Bernardis et al. 2009 (Opt. WL)

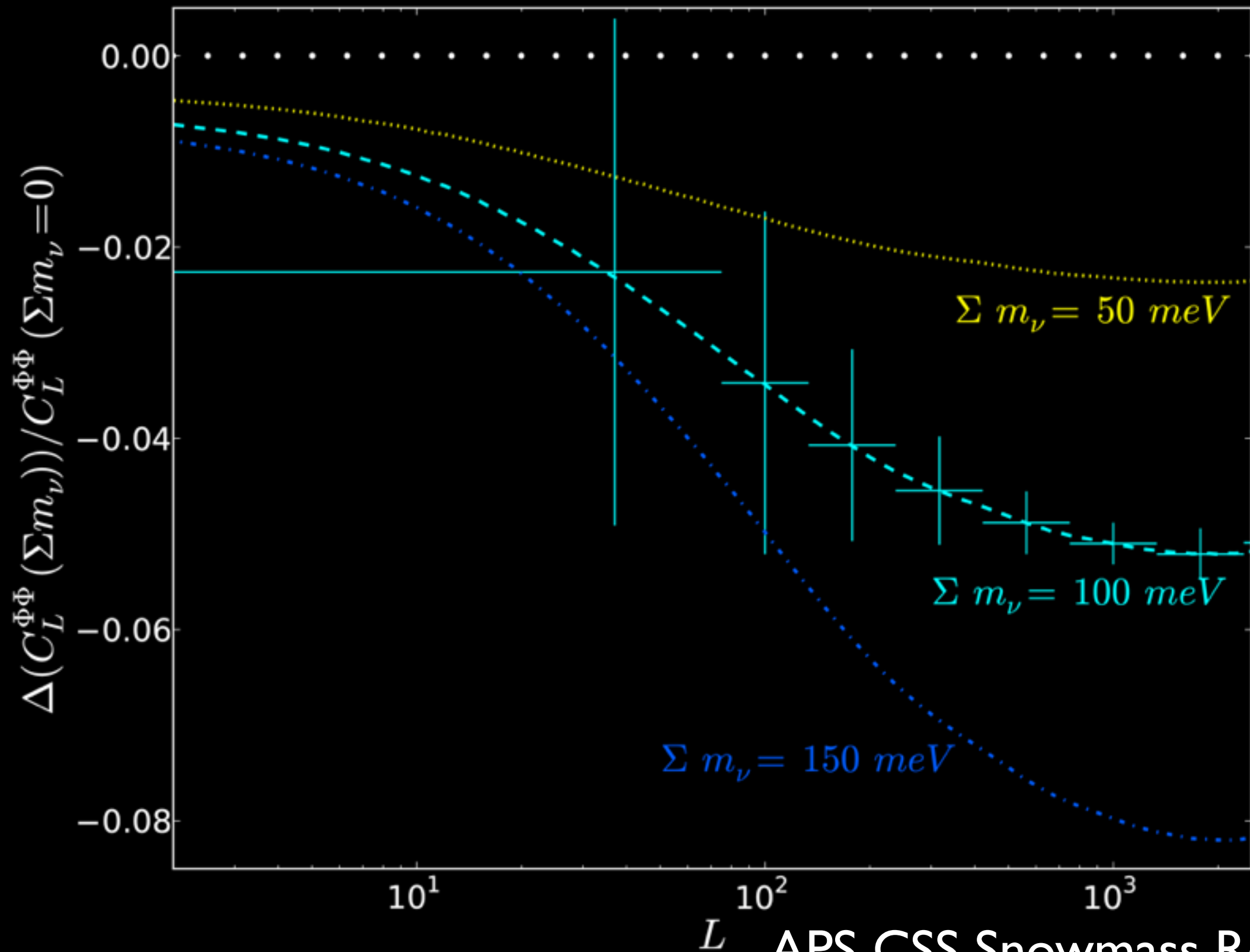
Joudaki & Kaplinghat 2011 (LSST)

Basse et al. 2013 (Euclid)

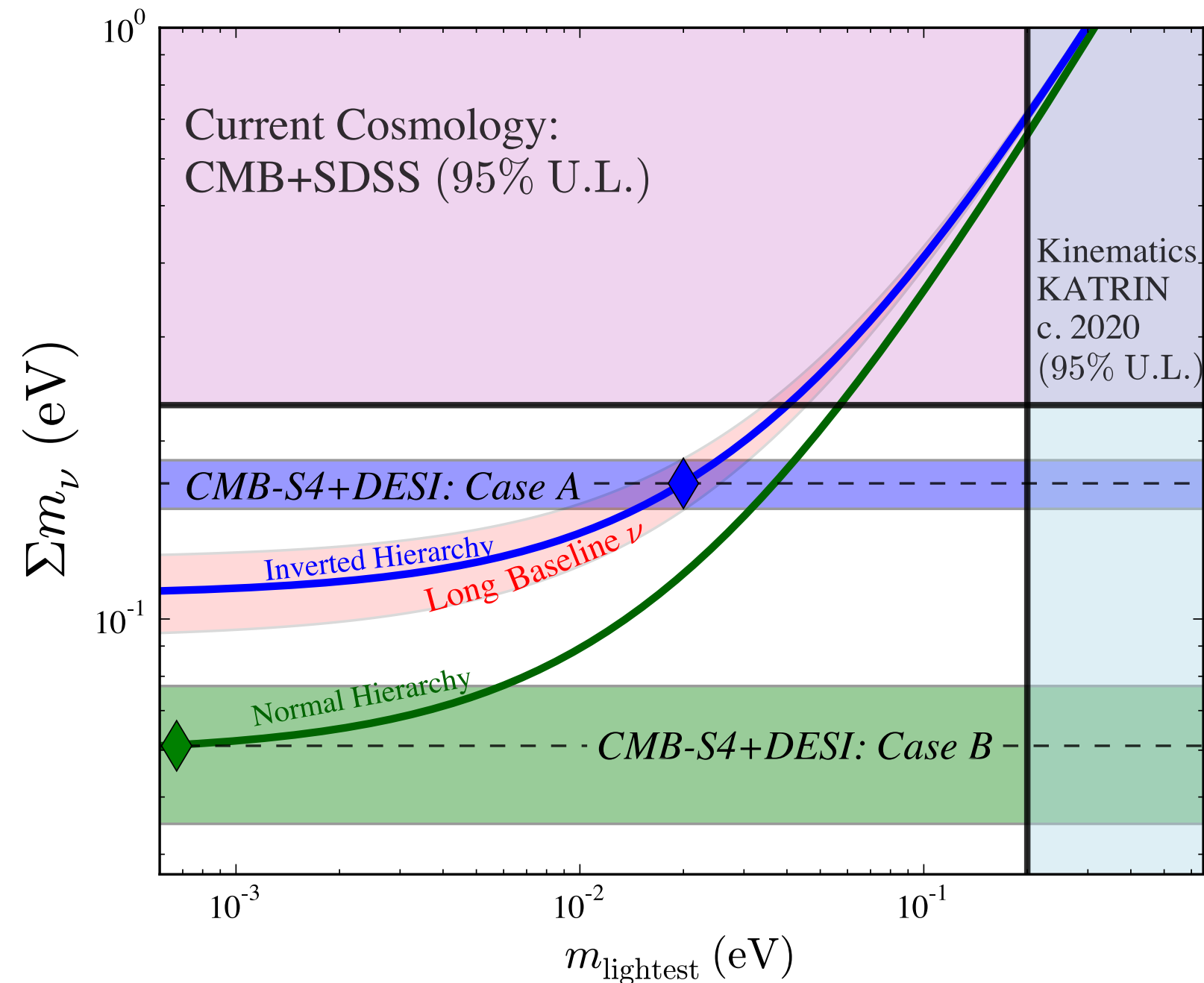
Abazajian et al. 2014 (Snowmass Report)

Wu et al. 2014 (CMB-S4 + DESI)

Lensing Potential Power: Relative Change over an Integrated $P(k)$



Cosmological & Laboratory Complementarity



CMB-S4 and DESI galaxy
survey BAO 1- σ
sensitivities are, at optimal
design:

$$\sigma(\Sigma m_\nu) = 15 \text{ meV}$$

$$\sigma(N_{\text{eff}}) = 0.016$$

providing $> 3\sigma$ sensitivity
to the oscillation-required
 $\Sigma m_\nu = 58 \text{ meV}$ and $> 2\sigma$
sensitivity to N_{eff}

CMB alone:

$$\sigma(\Sigma m_\nu) = 34 \text{ meV}$$

$$\sigma(N_{\text{eff}}) = 0.016$$

(Wu+ 2014)

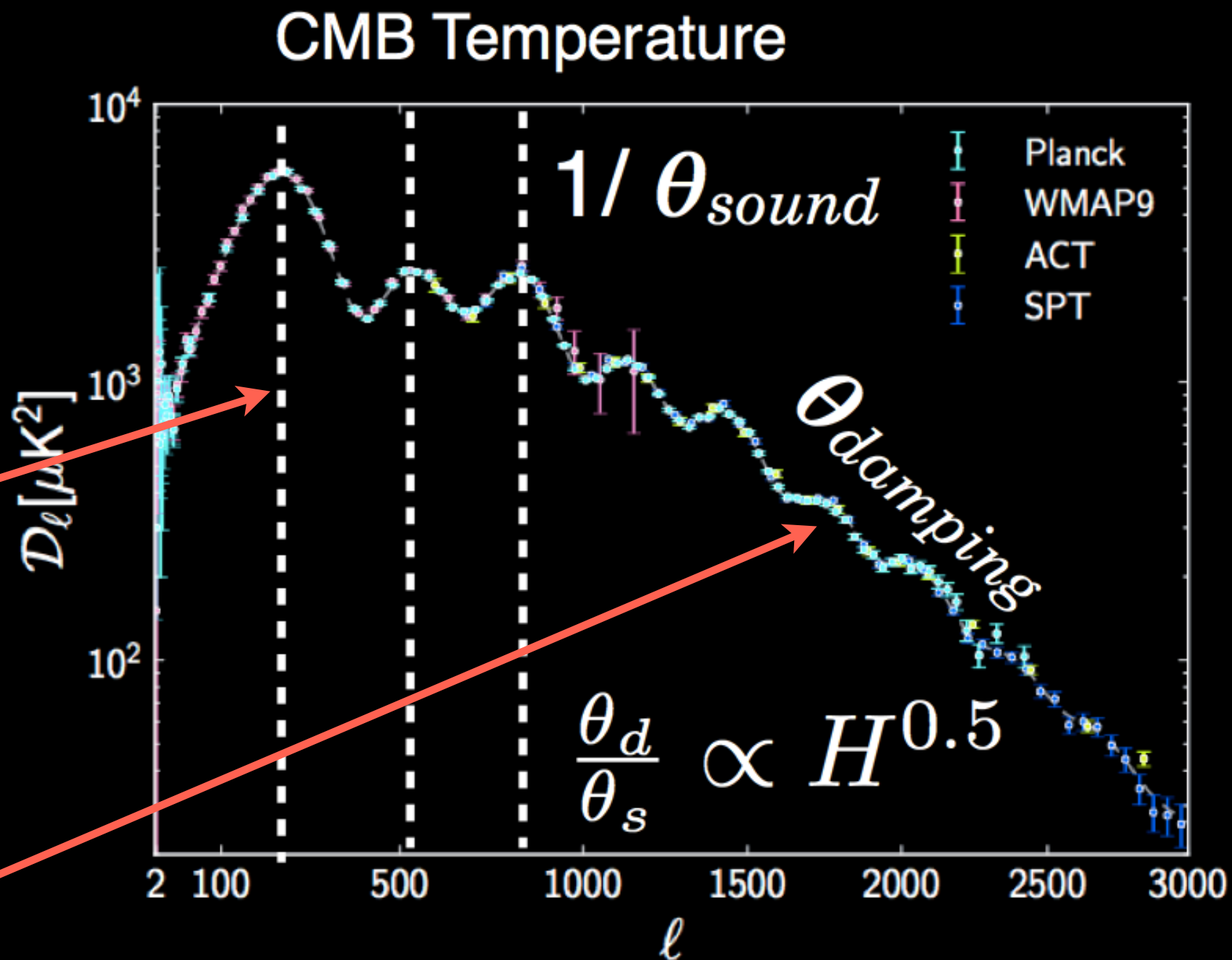
N_{eff} Effects on CMB

$$\frac{\theta_{\text{damping}}}{\theta_{\text{sound}}} \propto H^{1/2}$$

Larger N_{eff} Leads to
More Damping

Angular scale of
acoustic peaks
 $\theta_s \sim r_s/D$ is known
precisely

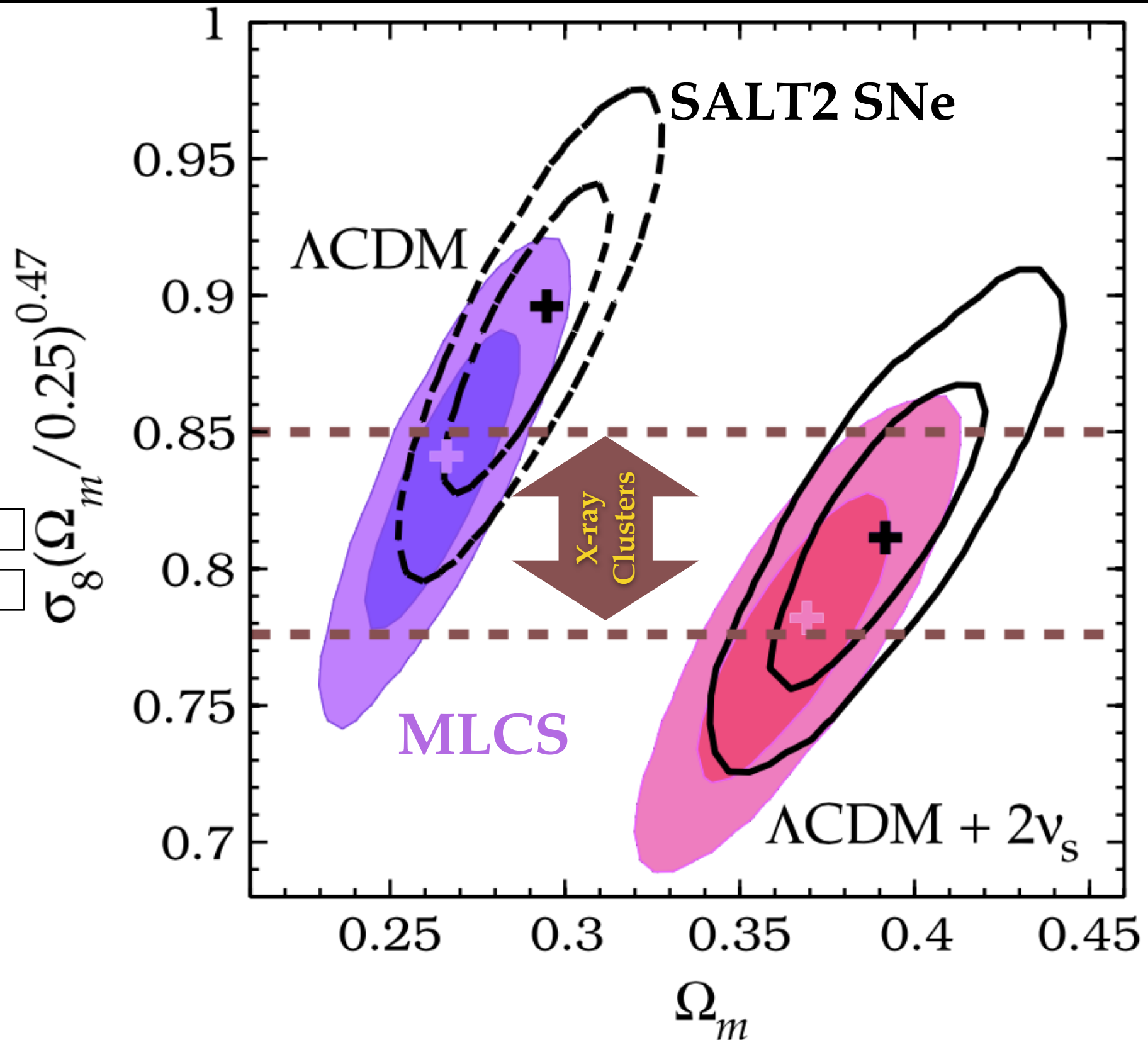
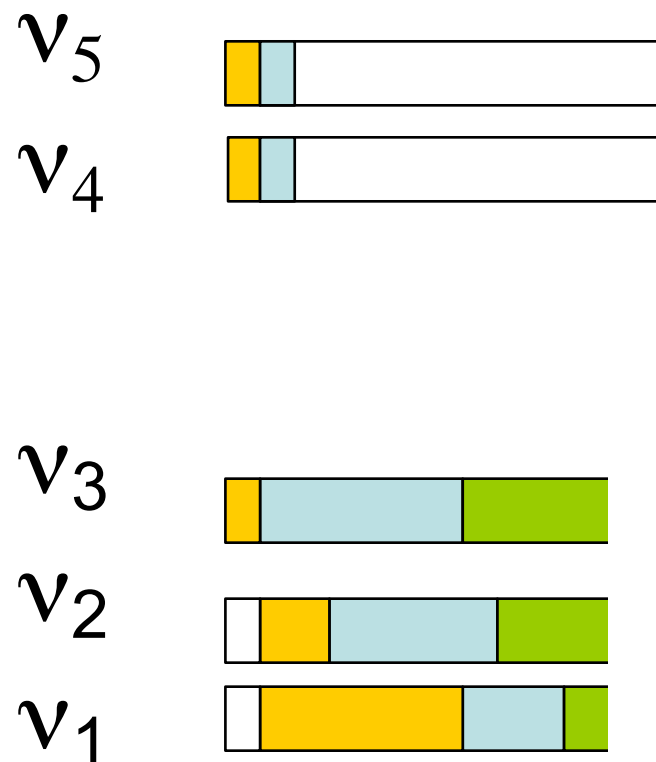
Angular scale of
damping $\theta_d \sim r_d/D$
measured recently



The Need for Cross-Analysis Between Neutrino Experiment and Cosmology

Joudaki, Abazajian & Kaplinghat 2013

⇒ Short baseline experiments require, minimally, 2 sterile ν 's with specific masses, mixings



Neutrino energy density beyond the standard model e.g., via lepton number generation

1. Extra Neutrino Flavors: $\nu_s, \nu_{s'}, \nu_{s''} \dots$

Models: 2+2, 3+1, 3+2, 3+N ...

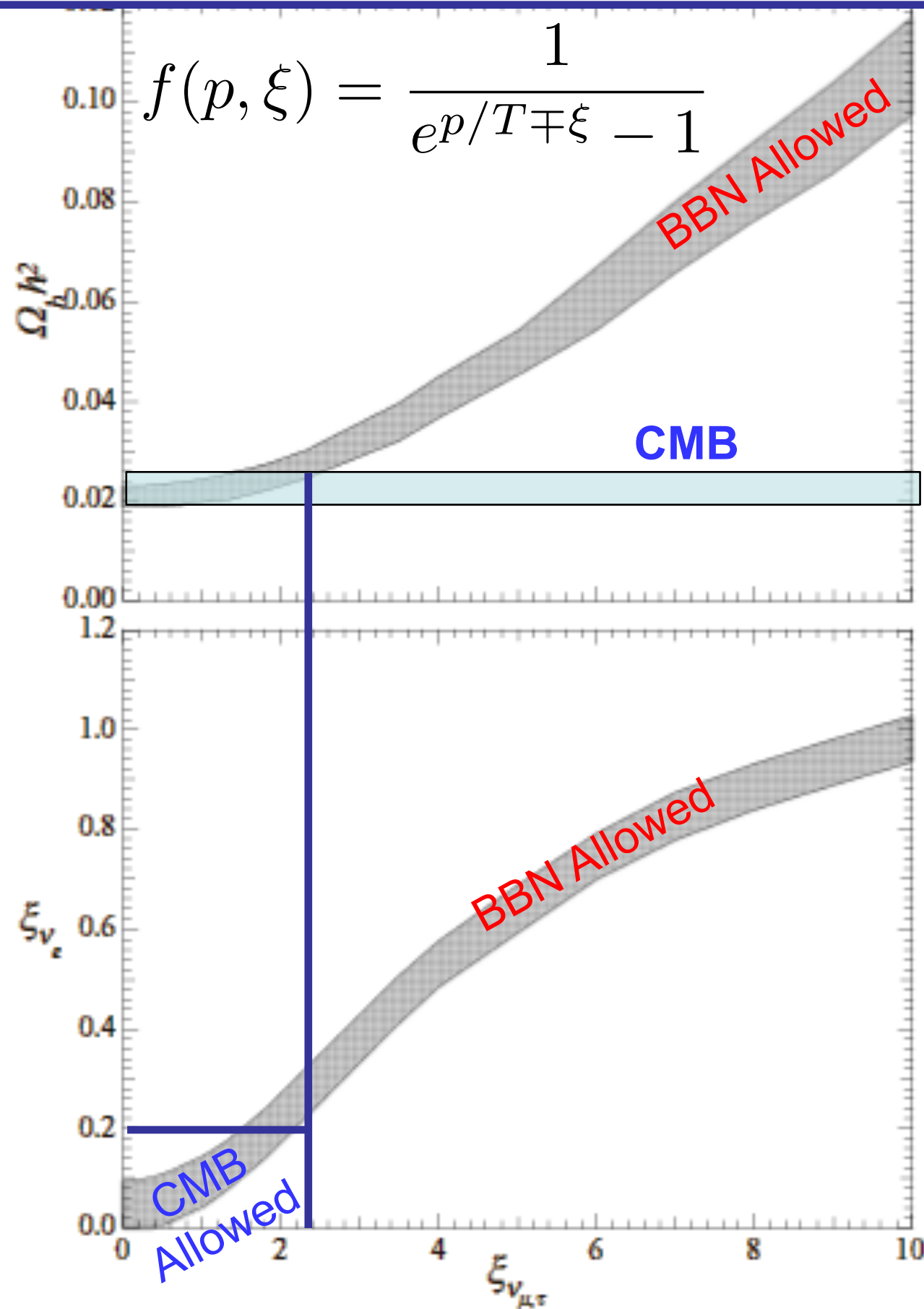
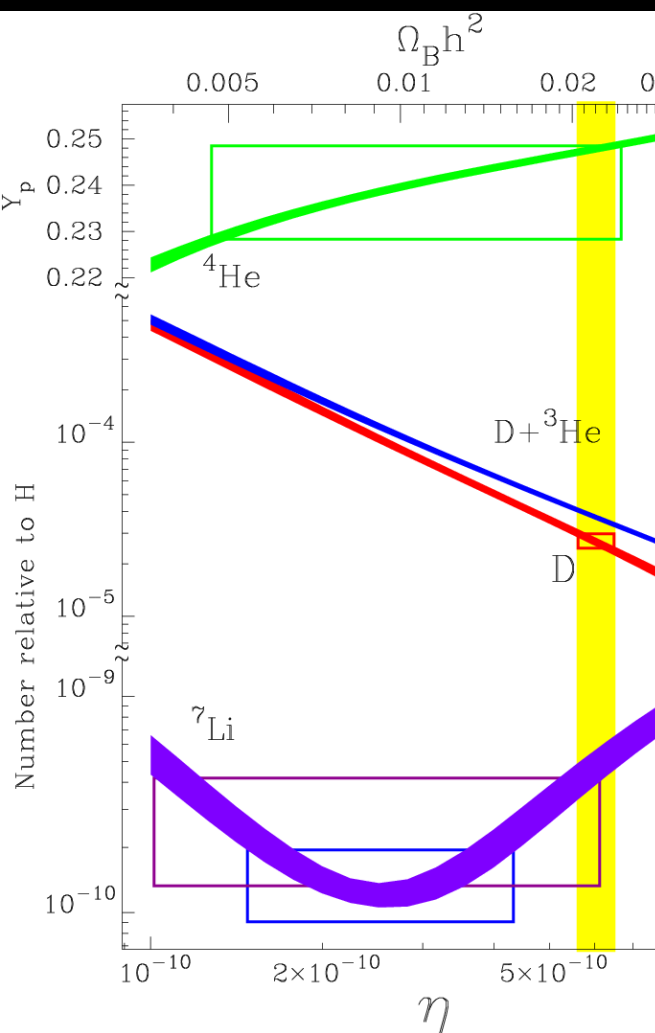
2. Neutrino Asymmetry from a leptogenesis:

$$\rho_\nu + \rho_{\bar{\nu}} = \frac{1}{2\pi^2} \int_0^\infty dp \, p^2 E_\nu (f_\nu(p) + f_{\bar{\nu}}(p))$$

$$f(p, \xi \equiv \mu/T) = \frac{1}{e^{p/T \mp \xi} - 1}$$

$$\Delta N_\nu = \frac{30}{7} \left(\frac{\xi}{\pi} \right)^2 + \frac{15}{7} \left(\frac{\xi}{\pi} \right)^4$$

“Degenerate” Big Bang Nucleosynthesis



Using the
observed
abundances of
 D , ^4He , and ^7Li
alone,

BBN is very
pliable to allow
large neutrino
asymmetries

Due to cancelling
effects between

$\nu_{\mu/\tau}$ and ν_e ,
and baryon
density, $\Omega_b h^2$

Orito et al '02

Quantum statistics with early universe neutrinos

Bloch Equations

$$\begin{aligned}\rho(p) &= \frac{1}{2} [P_0(p) + \mathbf{P}(p) \cdot \boldsymbol{\sigma}] \\ &= \frac{1}{2} \begin{pmatrix} P_0(p) + P_z(p) & P_x(p) - iP_y(p) \\ P_x(p) + iP_y(p) & P_0(p) - P_z(p) \end{pmatrix}\end{aligned}$$

Such that:

$$n_{\nu_\alpha}(p) = \frac{1}{2} [P_0(p) + P_z(p)] \quad n_{\nu_\beta}(p) = \frac{1}{2} [P_0(p) - P_z(p)]$$

Evolution of the density matrix for a homogeneous two-flavor neutrino gas plus background ($e^\pm, \mu^\pm \dots$) is generally

$$\dot{\mathbf{P}}(p) = \mathbf{V}(p) \times \mathbf{P}(p) + [R_\alpha(p) - R_\beta(k)] \hat{\mathbf{z}} + \mathcal{C}[\mathbf{P}(p)]$$

$$\dot{P}_0(p) = R_\alpha(p) + R_\beta(p)$$

Coherent Behavior

$$\partial_t \mathbf{P}(p) = \mathbf{V}(p) \times \mathbf{P}(p)$$

$$\mathbf{V}(p) = \Delta(p) + [V^B(p) + V^T(p)] \hat{\mathbf{z}} + V^S(p)$$

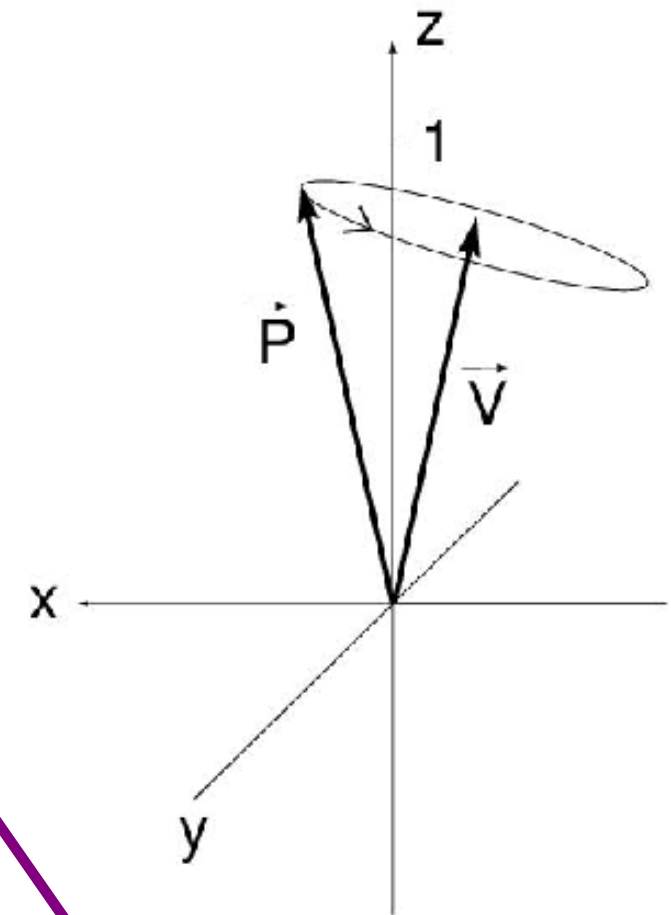
...dipole precession in a magnetic field

Vacuum Oscillation

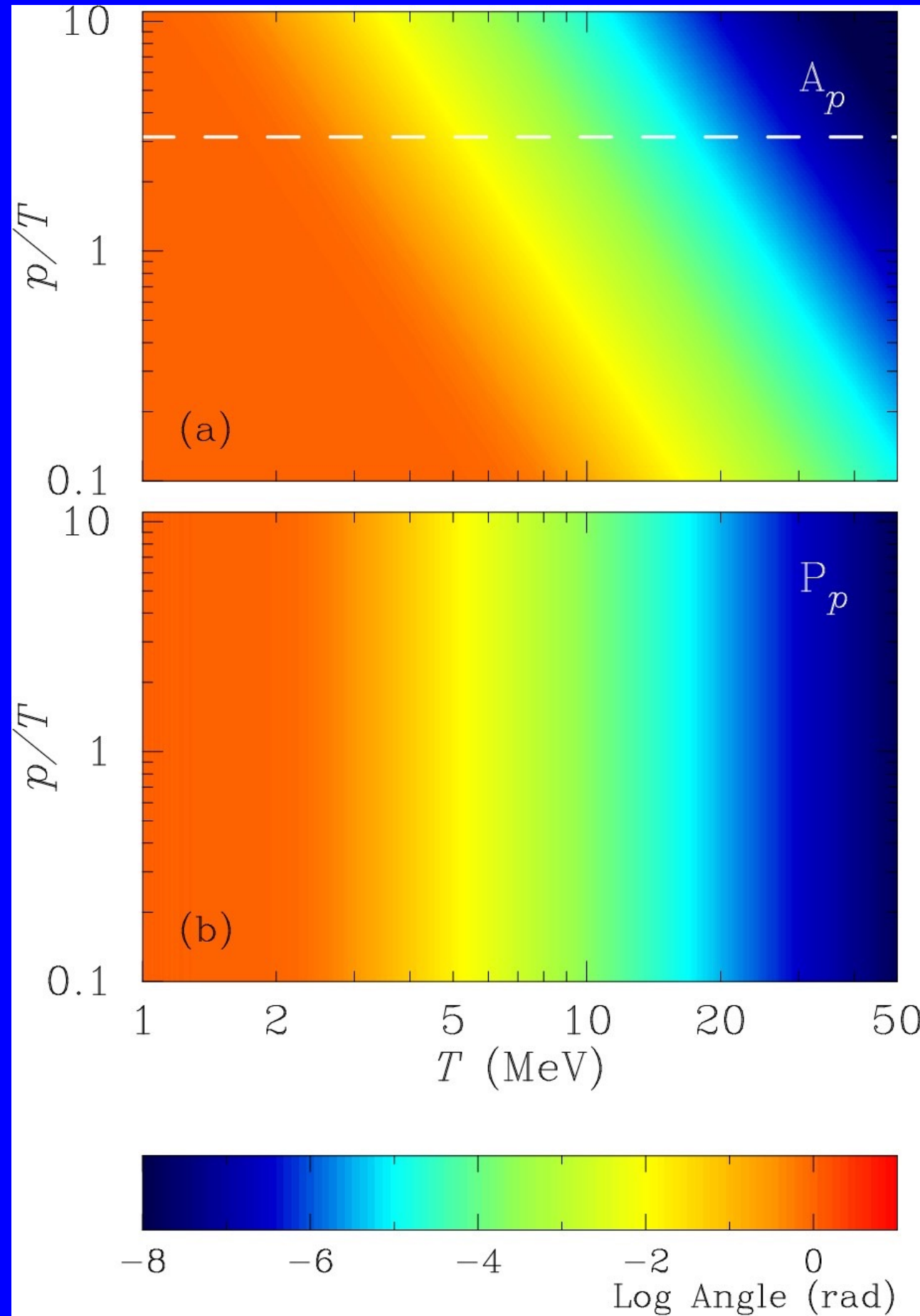
$e^{+/-} / \mu^{+/-}$ background

**Thermal Potential
(finite temperature effects)**

Neutrino Self-Potential



Neutrino Flavor Momentum Synchronization...



Collectively:

$$\partial_t \mathbf{I} = \mathbf{A}_{\text{eff}} \times \mathbf{I}$$

Assuming $\mathbf{I} \parallel (\mathbf{P}_p + \bar{\mathbf{P}}_p)$:

$$\mathbf{A}_{\text{eff}} \simeq \frac{1}{I^2} \int \mathbf{A}_p (\mathbf{P}_p + \bar{\mathbf{P}}_p) \, y$$

follows...

$$\mathbf{A}_{\text{eff}} \equiv \Delta_{\text{sync}} (\sin 2\theta_{\text{sync}} \hat{\mathbf{x}} - \cos 2\theta_{\text{sync}} \hat{\mathbf{z}}),$$

$$\frac{p_{\text{sync}}}{T} = \pi \sqrt{1 + \xi^2 / 2\pi^2} \simeq \pi$$

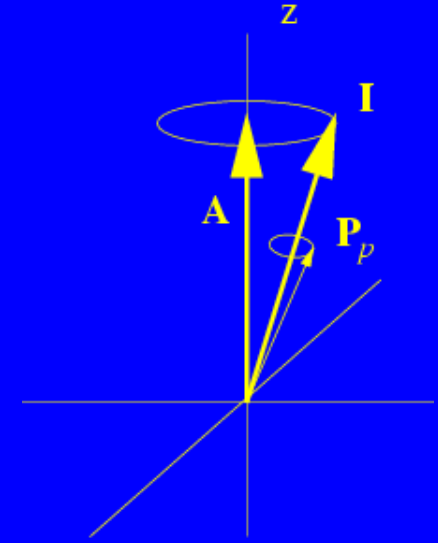
while the average for a FD distribution is

$$\langle p/T \rangle = \frac{7\pi^4}{180\zeta(3)} \simeq 3.15$$

$$\Delta(p) = \frac{\delta m^2}{2p} (\sin 2\theta_0, 0, -\cos 2\theta_0)$$

Large to maximal mixing is essential...

Abazajian, Beacom & Bell
2002





The Death of Degenerate BBN and New Constraints...

Constraints on ν_e asymmetry:

$$\Delta\xi_e \approx \frac{\Delta Y_p}{Y_p}$$

Conservatively, errors

$$\frac{\Delta Y_p}{Y_p} \approx 0.1 \Rightarrow \Delta\xi_e$$

or, more accurate accounting gives

$$0.03 \lesssim \xi_e \lesssim 0.04$$

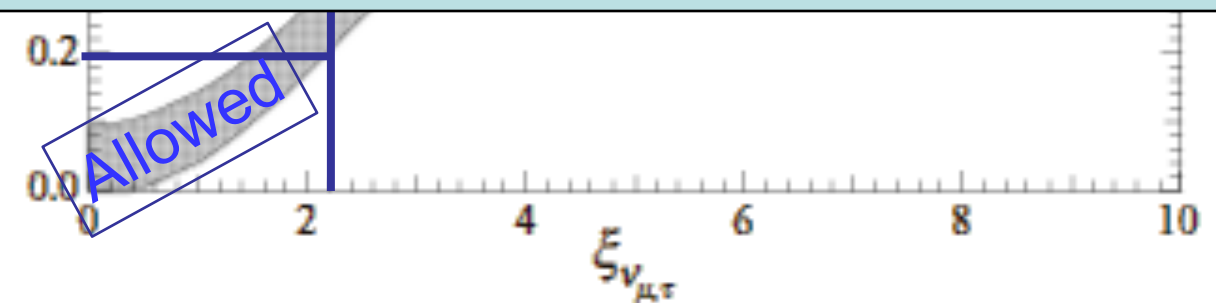
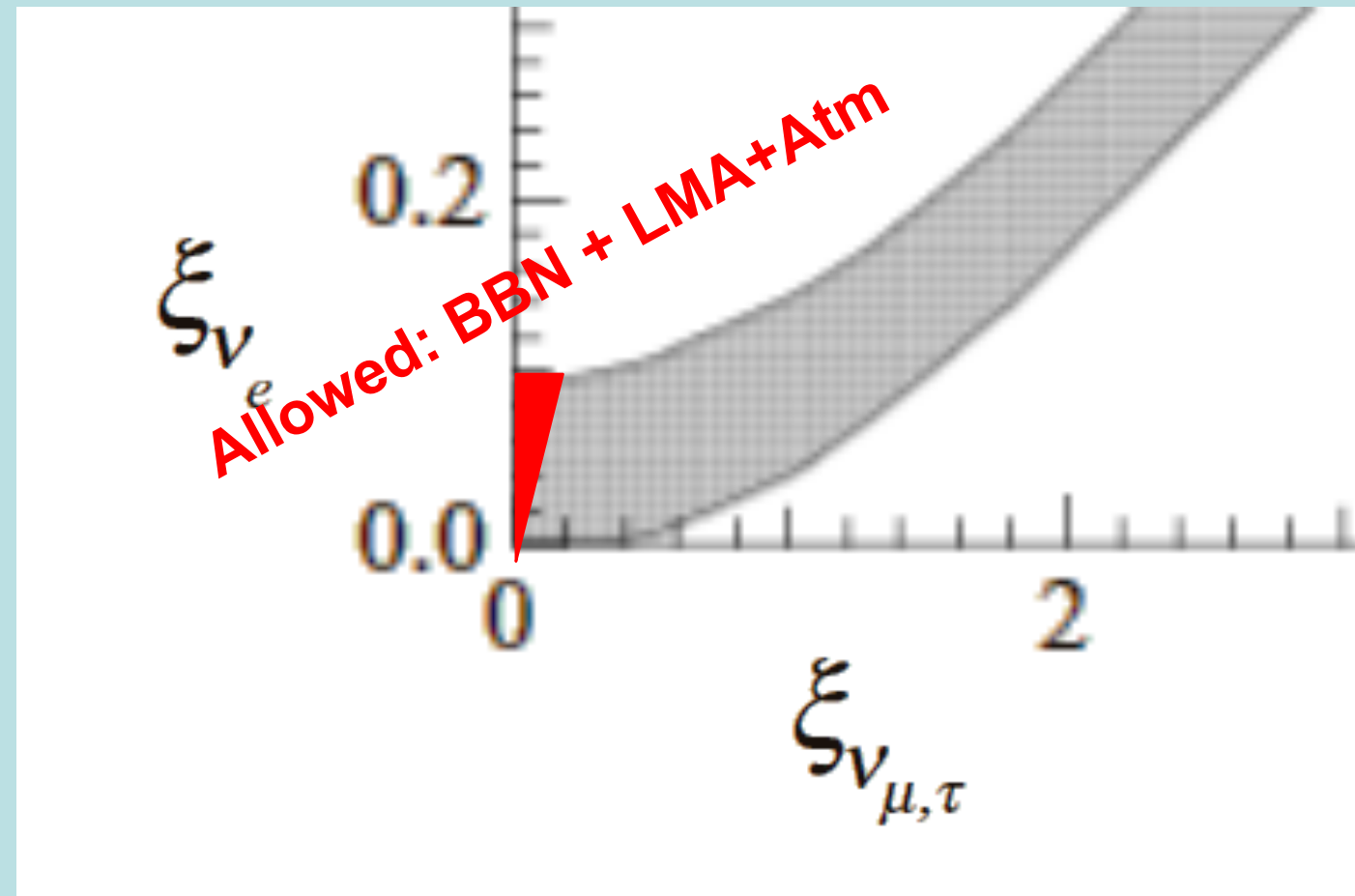
$$|L_e| \lesssim 0.1$$

and

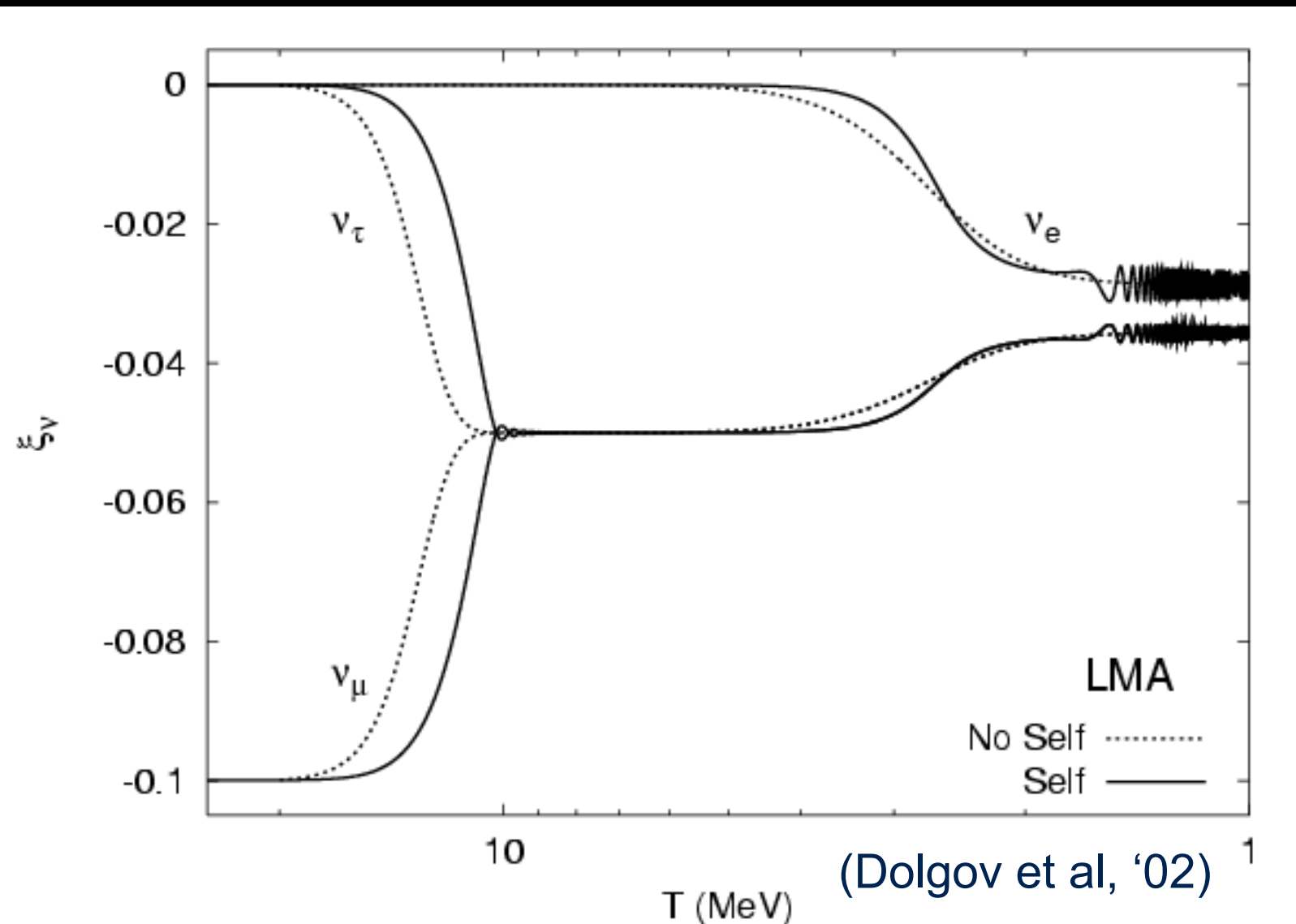
$$\xi_e^f = \left(\frac{1 - \cos 2\theta_0}{2} \right)$$

$$\xi_{\mu^*}^f = \left(\frac{1 + \cos 2\theta_0}{2} \right)$$

$$\Rightarrow |L_\mu + L_\tau| \lesssim 0.5$$



Transformation of neutrino asymmetries



LMA solar + maximal ATM

→ Any chemical potential will alter the ^4He abundance.

$$\xi \ll 1 \implies \Delta N_\nu < 0.04$$

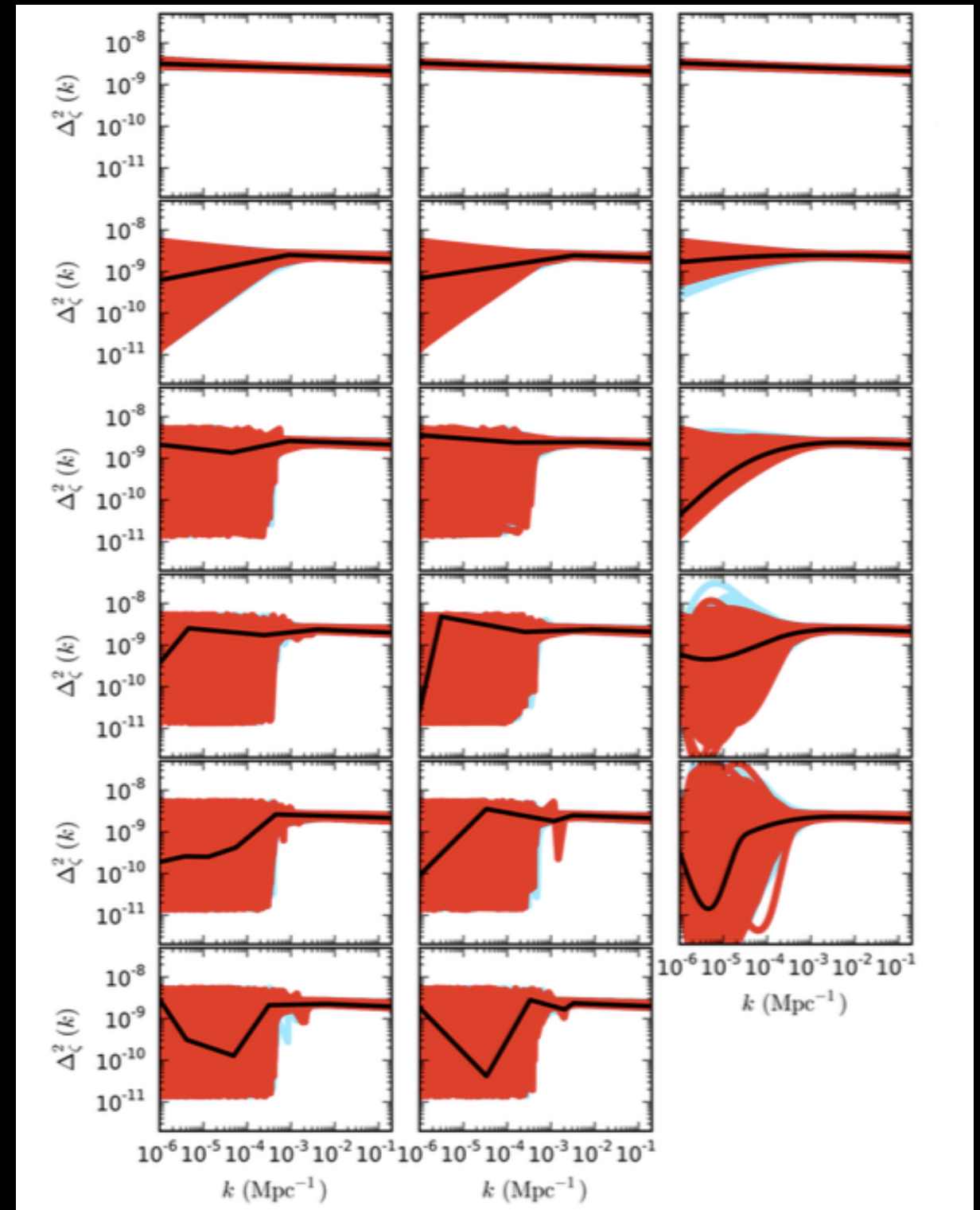
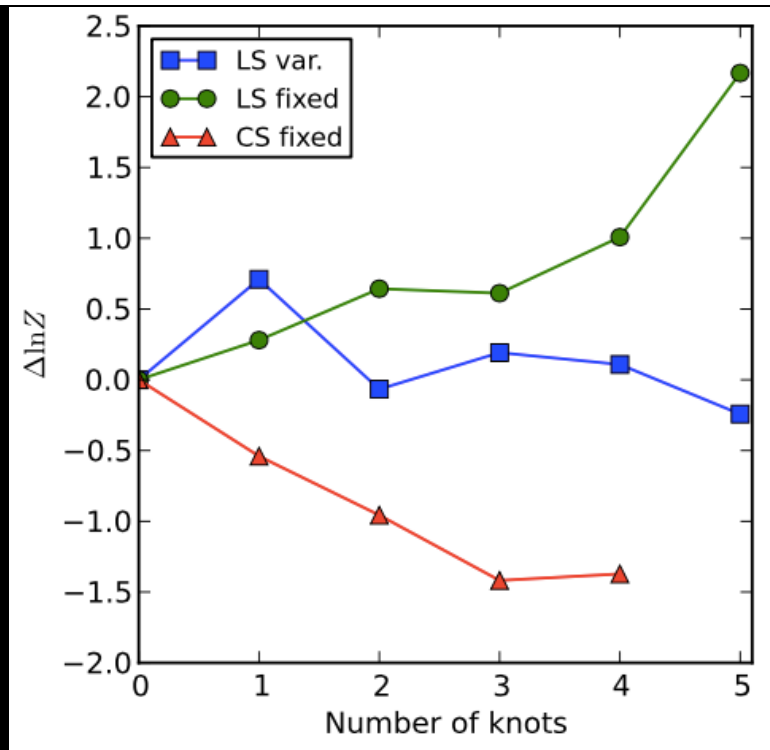
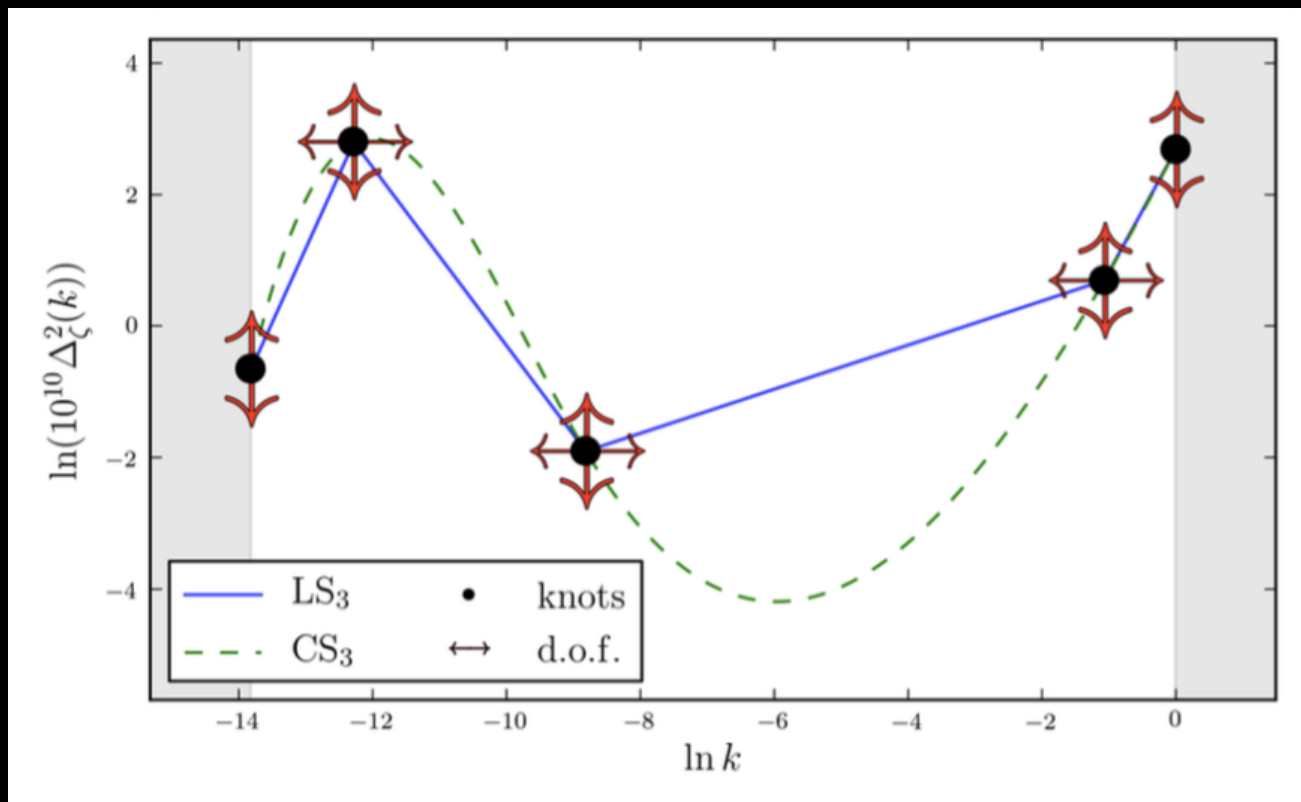
The claimed standard abundance of $N_{\text{eff}}(\text{standard}) = 3.046$ is uncertain to within this order of magnitude due to oscillation effects during decoupling

More general density:

$$n_\nu = \left(\frac{3}{11} \right) \frac{2\zeta(3)}{\pi^2} T^3 F_2(\xi) \\ \approx 113 \text{ cm}^{-3} F_2(\xi)$$

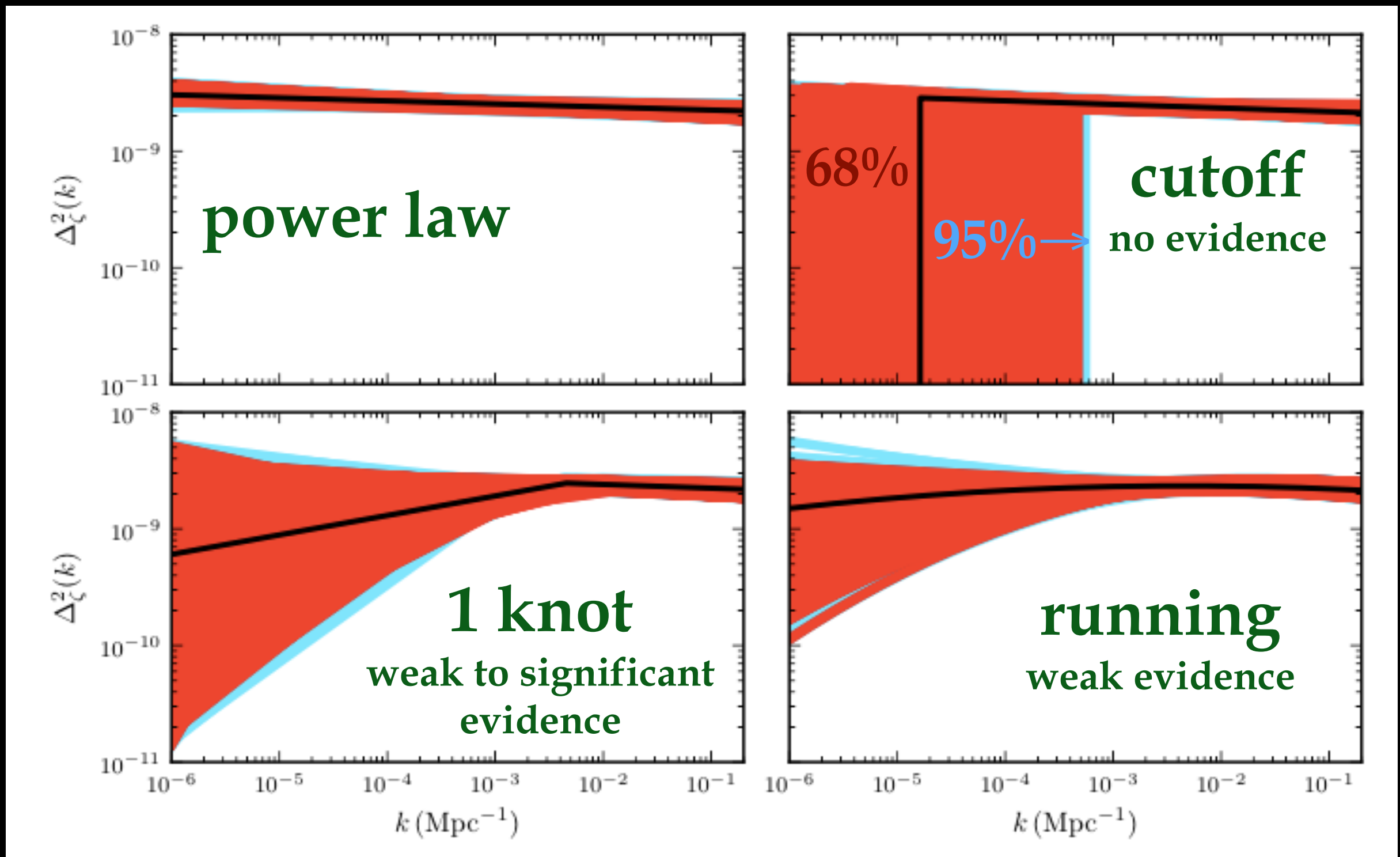
Signals of deviations from a power law primordial power spectrum in the CMB?

Aslanyan, Price, Abazajian & Easter 2014



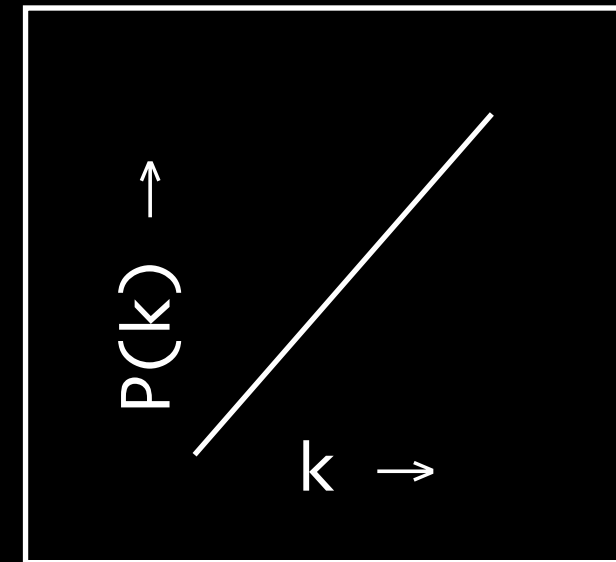
Signals of deviations from a power law PPS in the CMB **with BICEP2 relic *B*-modes?**

Abazajian, Aslanyan, Easter & Price 2014

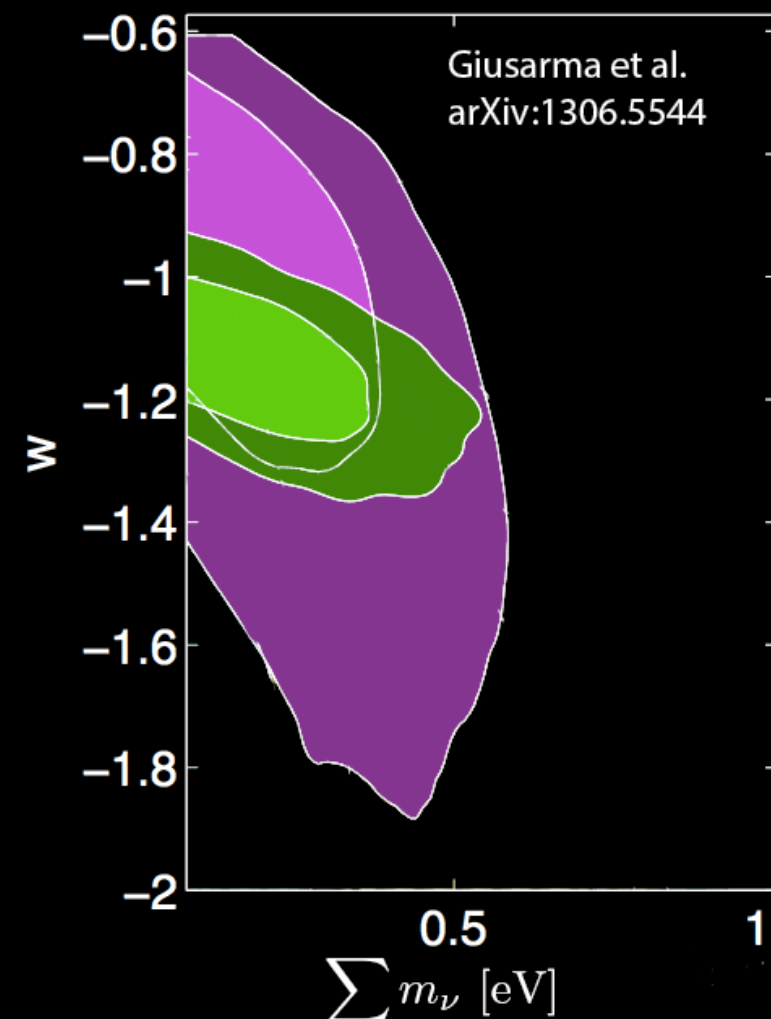


Neutrino Mass from Cosmology: What would break if cosmology and neutrino experiment disagree?

1. Primordial power spectrum $P(k)$ is a simple power law



2. No other prevalent “non-vanilla” cosmological parameters and physics: w , N_{eff} , modified gravity...



Summary

- Cosmology has the strongest inferred experimental sensitivity on the total neutrino mass, and are forecast to maintain that position.
- CMB-S4 and DESI galaxy survey 1- σ sensitivities are:

$$\sigma(\Sigma m_\nu) = 15 \text{ meV} \text{ \& } \sigma(N_{\text{eff}}) = 0.016$$

providing $> 3\sigma$ sensitivity to the oscillation-required $\Sigma m_\nu = 58 \text{ meV}$ and $> 2\sigma$ sensitivity to N_{eff} (Wu+ 2014)

- The neutrino background is not necessarily the simplest possible case. Asymmetries from lepton number generation, extra neutrino density from massive particle decay, and new neutrino interactions could all reveal themselves.
- *Other relativistic energy density can be present.*
- **What if we do not detect the minimal model?**
*If the minimal neutrino sector, with $\Sigma m_\nu = 58 \text{ meV}$ and $N_{\text{eff}} = 3.046$, is **not** robustly detected, it would imply something is “off” in another aspect or aspects of cosmology, including possibly: non-constant dark energy, a non-power-law primordial perturbation spectrum, extra particle or radiation species, non-zero curvature, as well as other possibilities, e.g., a nonthermal cosmological neutrino background.*